New approach to the local strong parity violation in the quark-gluon plasma



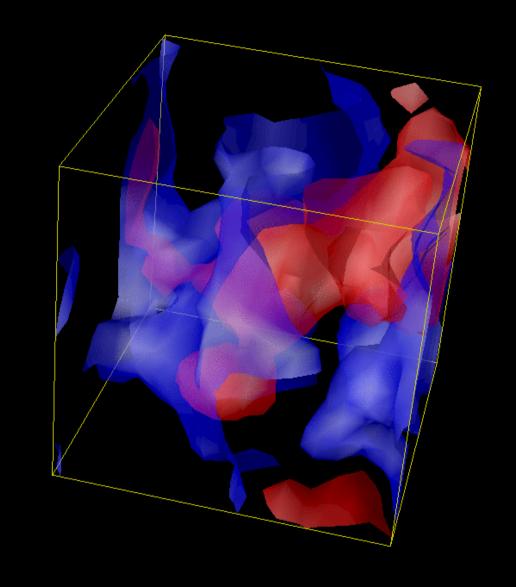
Tigran Kalaydzhyan

QCD vacuum

 $G^{a\mu
u} ilde{G}^a_{\mu
u}$

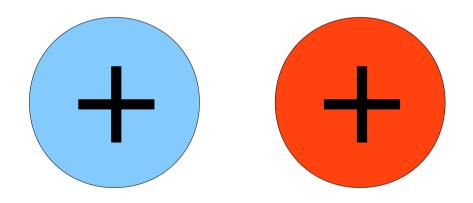


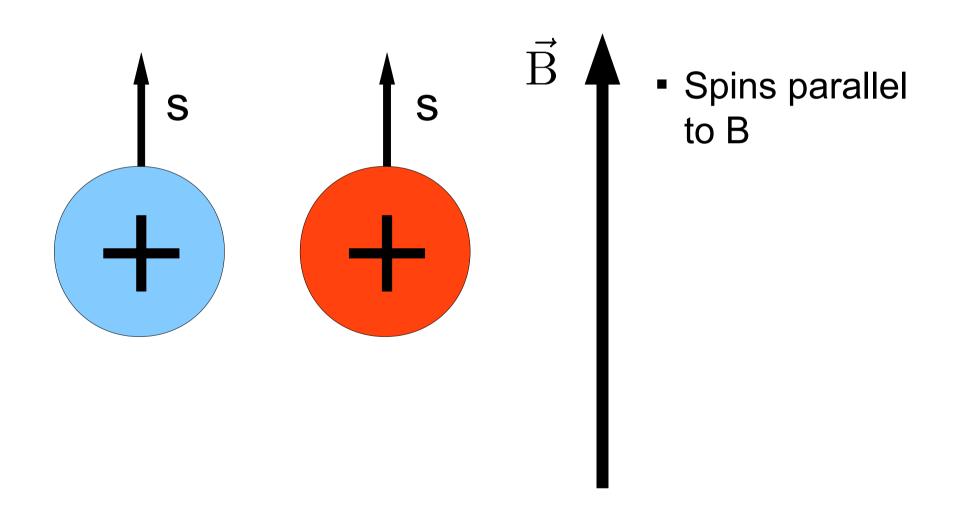
$$\rho_R \neq \rho_L$$



Positive topological charge density

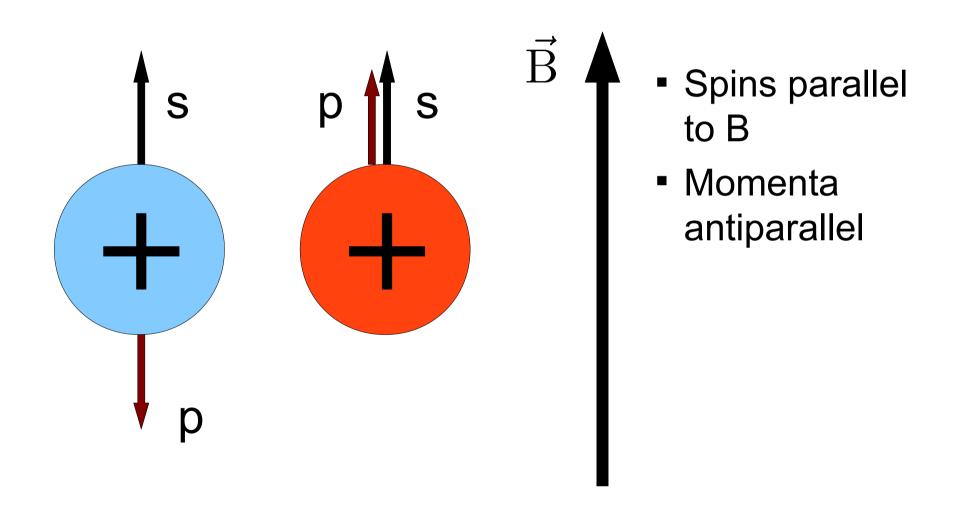
Negative topological charge density





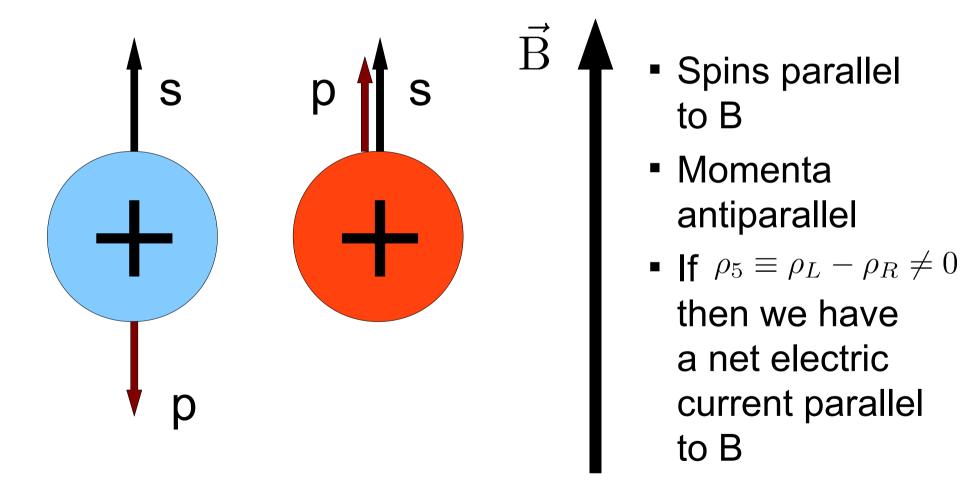
Right-handed

Left-handed



Right-handed

Left-handed

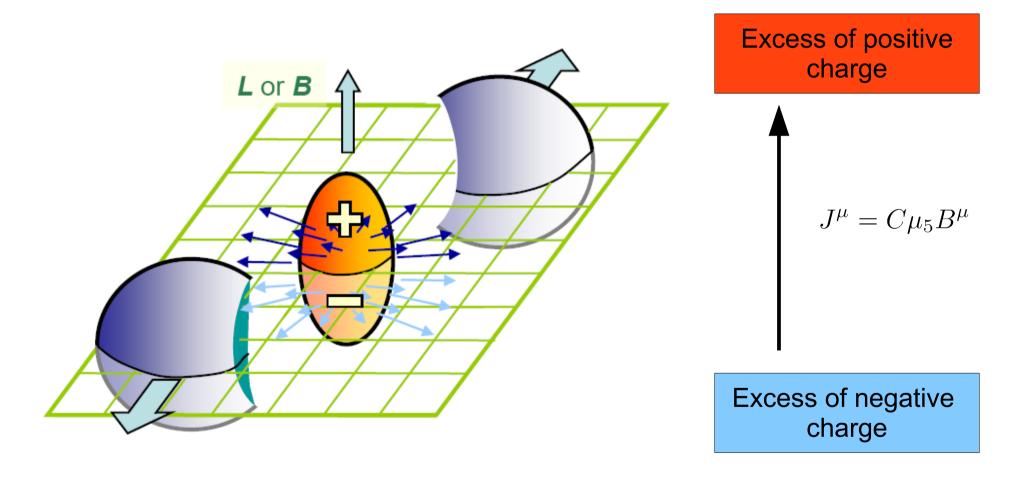


Left-handed

Right-handed

Kharzeev, McLerran, Warringa (2007)

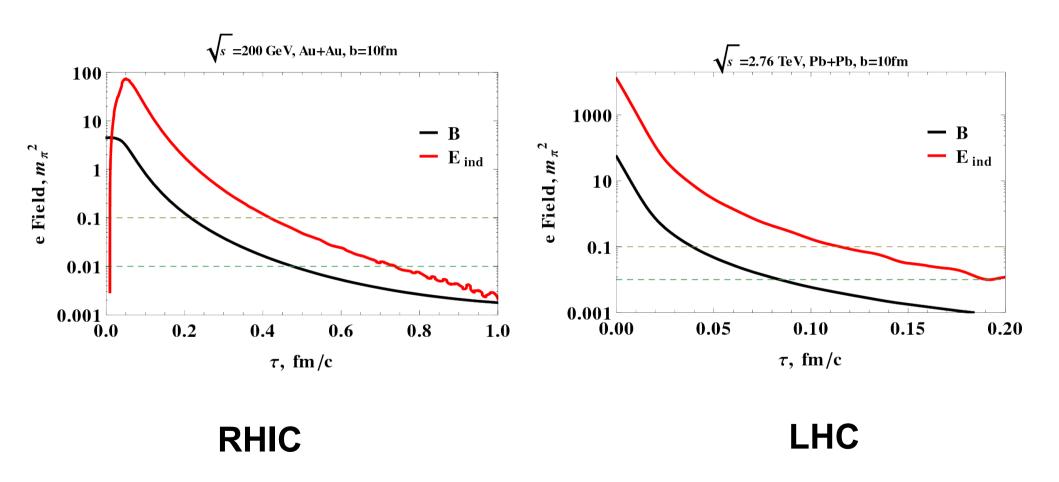
Chiral Magnetic Effect



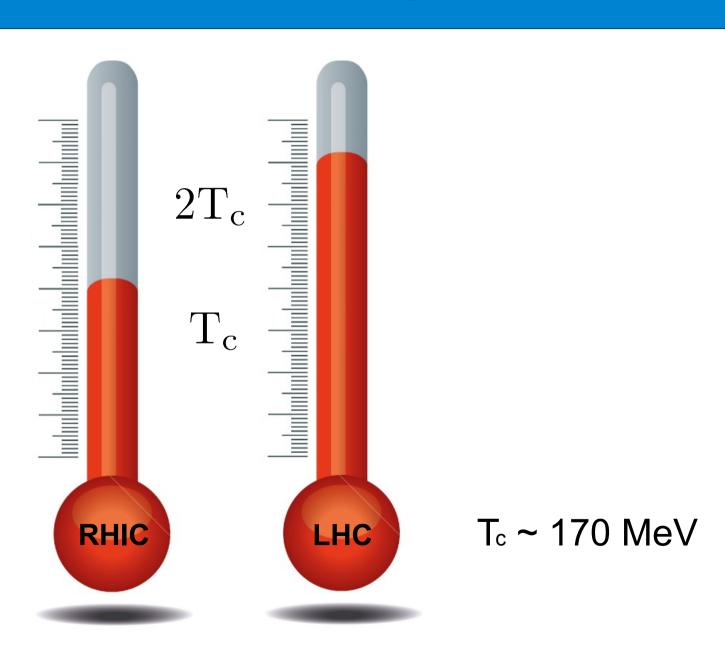
Fukushima, Kharzeev, McLerran, Warringa (2007)

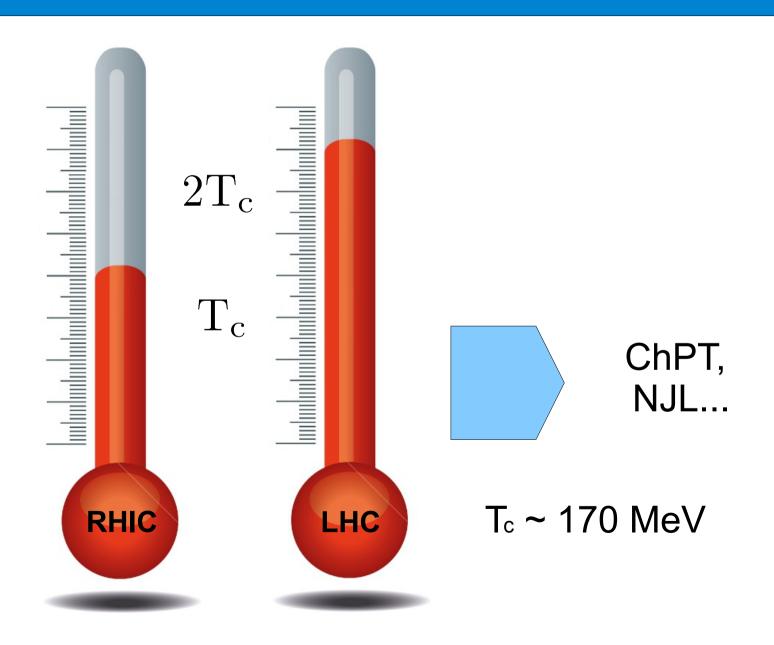
For a local strong parity violation see e.g. 0909.1717 (STAR) and 1207.0900 (ALICE)

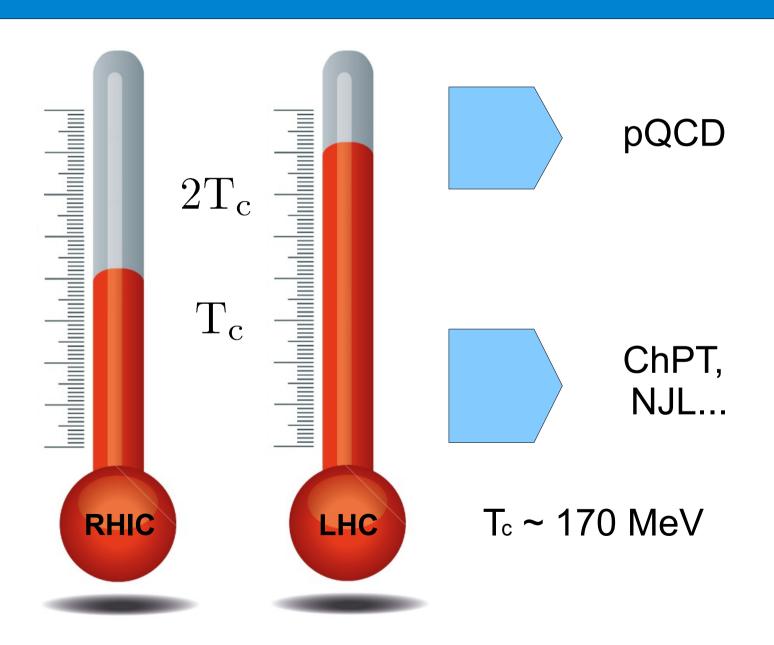
Electromagnetic fields

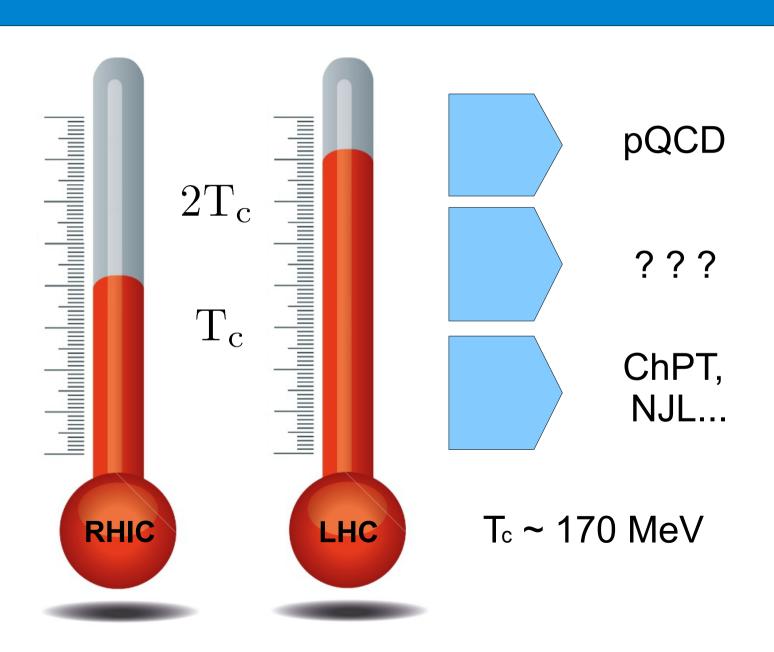


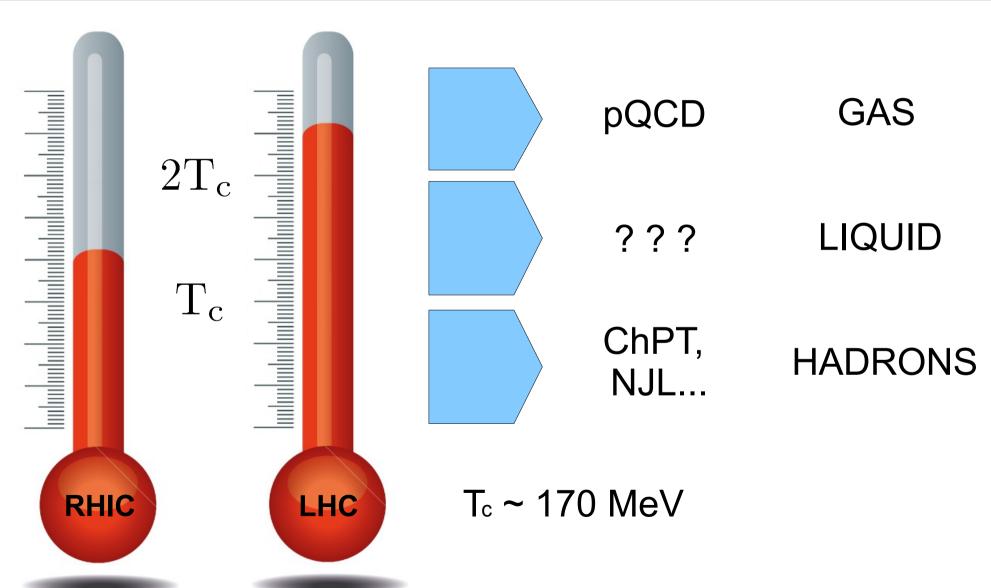
Huge electromagnetic fields, never observed before!











Shuryak (2003), Chernodub and Zakharov (2008)

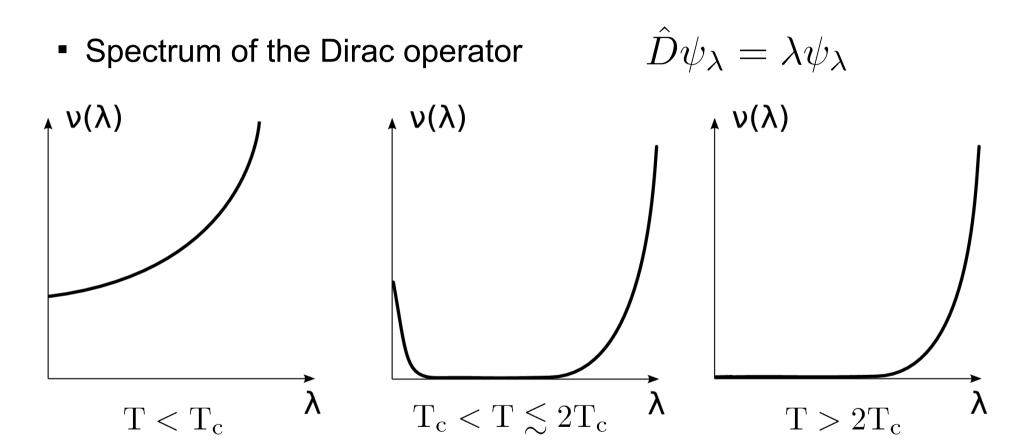
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- Solve the hydro equations in gradient expansion, satisfying thermodynamic laws
- Extract phenomenological output for the heavy-ion collisions

Insight from the lattice

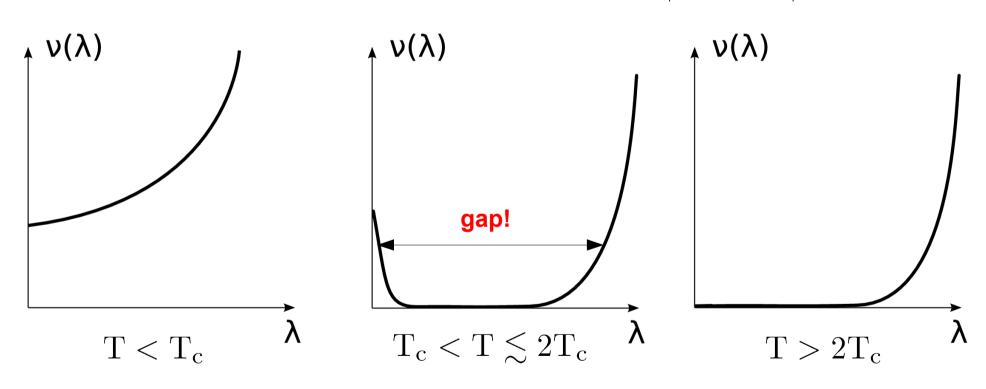


Chiral properties are described by near-zero modes

Insight from the lattice

Spectrum of the Dirac operator

$$\hat{D}\psi_{\lambda} = \lambda\psi_{\lambda}$$



- Chiral properties are described by near-zero modes
- There are two separated parts of the spectrum at intermediate temperatures!

Why "superfluidity"?

Energy Normal motions **Curl-free motions**

AUGUST 15, 1941

PHYSICAL REVIEW

Theory of the Superfluidity of Helium II

L. LANDAU
Institute of Physical Problems, Academy of Sciences USSR, Moscow, USSR

Therefore, between the lowest energy levels of vortex and potential motion there must be a certain energy interval Δ .

The supposition that the normal level of potential motions lies lower than the beginning of the spectrum of vortex motions leads to the phenomenon of superfluidity.

One of these motions is "normal" and the other is "superfluid."

We will not consider any spontaneously broken symmetry!

• Euclidean functional integral for $SU(N_c) \times U_{em}(1)$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(D\!\!\!\!/ - im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

$$D = -i(\partial + A + g G + \gamma_5 A_5).$$

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where we define the Dirac operator as

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integrate out quarks below a cut-off Dirac eigenvalue Λ.

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- consider a pure gauge $A_{5\mu}=\partial_{\mu}\theta$ for the auxiliary axial field
- and the chiral limit $m \to 0$

The total effective Euclidean Lagrangian reads as

$$\mathcal{L}_{E}^{(4)} = \frac{1}{4} G^{a\mu\nu} G^{a}_{\mu\nu} + \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - j^{\mu} A_{\mu}$$

$$+ \frac{\Lambda^{2} N_{c}}{4\pi^{2}} \partial^{\mu} \theta \partial_{\mu} \theta + \frac{g^{2}}{16\pi^{2}} \theta G^{a\mu\nu} \widetilde{G}^{a}_{\mu\nu} + \frac{N_{c}}{8\pi^{2}} \theta F^{\mu\nu} \widetilde{F}_{\mu\nu}$$

$$+ \frac{N_{c}}{24\pi^{2}} \theta \Box^{2} \theta - \frac{N_{c}}{12\pi^{2}} \left(\partial^{\mu} \theta \partial_{\mu} \theta \right)^{2}$$

So we get an axion-like field with decay constant $f=\frac{2\Lambda}{\pi}\sqrt{N_c}$ and a negligible mass $m_{\theta}^2=\lim_{V\to\infty}\frac{\langle Q^2\rangle}{f^2V}\equiv\chi(T)/f^2$.

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Dynamical axion-like internal degree of freedom in QCD!

Interpretation of the scale \(\Lambda\)

• From the quartic Lagrangian at $N_c=N_f=1$ we get

$$\rho_5 = -\lim_{t_E \to 0} \frac{\delta \mathcal{L}_E^{(4)}}{\delta \mu_5} = \frac{1}{2} \left(\frac{\Lambda}{\pi}\right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

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- Free quarks and a strong B-field: $\Lambda = 2\sqrt{|eB|}$
- Dynamical fermions (1105.0385): $\Lambda \simeq 3\,\mathrm{GeV} \gg \Lambda_{QCD}$

One more remark

"Axionic" part of the Lagrangian

$$\mathcal{L}_{\theta} = \frac{\Lambda^2 N_c}{4\pi^2} \partial^{\mu}\theta \partial_{\mu}\theta + \frac{N_c}{24\pi^2} \theta \Box^2 \theta - \frac{N_c}{12\pi^2} \left(\partial^{\mu}\theta \partial_{\mu}\theta\right)^2 + \dots$$

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If magnetic field dominates over other scales, then we can make the following redefinition: $\theta \to \frac{\pi}{\sqrt{2N_c eB}}\theta$

$$\mathcal{L}_{\theta} \to \frac{1}{2} \partial^{\mu} \theta \partial_{\mu} \theta + \frac{1}{48 eB} \theta \Box^{2} \theta - \frac{\pi^{2}}{48 N_{c}(eB)^{2}} (\partial^{\mu} \theta \partial_{\mu} \theta)^{2} + \dots$$

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In the limit $B \to \infty$ bosonization becomes exact, which is an evidence of the $(3+1) \to (1+1)$ reduction!

Hydrodynamic equations

Considering EOM for the Minkowski effective Lagrangian and only the color-singlet states, we obtain:

$$\partial_{\mu}T^{\mu\nu} = F^{\nu\lambda}J_{\lambda} \,,$$

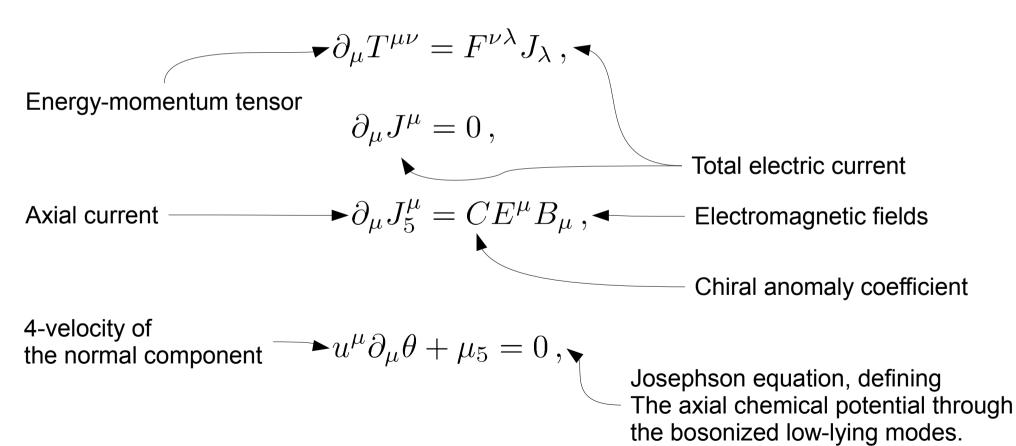
$$\partial_{\mu}J^{\mu}=0\,,$$

$$\partial_{\mu}J_{5}^{\mu} = CE^{\mu}B_{\mu} \,,$$

$$u^{\mu}\partial_{\mu}\theta + \mu_5 = 0,$$

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Similar to the superfluid dynamics!

Hydrodynamic equations

 Solving the hydrodynamic equations in the gradient expansion, we obtain the constitutive relations:

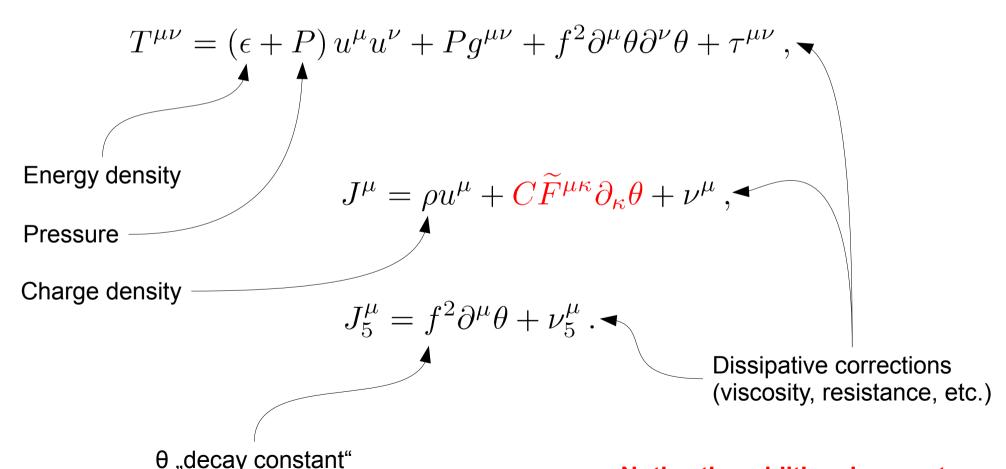
$$T^{\mu\nu} = (\epsilon + P) u^{\mu}u^{\nu} + Pg^{\mu\nu} + f^2 \partial^{\mu}\theta \partial^{\nu}\theta + \tau^{\mu\nu} ,$$

$$J^{\mu} = \rho u^{\mu} + C\widetilde{F}^{\mu\kappa} \partial_{\kappa} \theta + \nu^{\mu} ,$$

$$J_5^{\mu} = f^2 \partial^{\mu} \theta + \nu_5^{\mu} .$$

Hydrodynamic equations

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Notice the additional current

Higher order correction obey the Landau conditions

$$u_{\mu}\tau^{\mu\nu} = 0, \qquad u_{\mu}\nu^{\mu} = 0, \qquad u_{\mu}\nu^{\mu}_{5} = 0$$

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Using both hydrodynamic equations and constitutive relations one can derive

$$\partial_{\mu}(su^{\mu} - \frac{\mu}{T}\nu^{\mu} - \frac{\mu_{5}}{T}\nu^{\mu}_{5}) = -\frac{1}{T}(\partial_{\mu}u_{\nu})\tau^{\mu\nu} - \nu^{\mu}(\partial_{\mu}\frac{\mu}{T} - \frac{1}{T}E_{\mu}) - \nu^{\mu}_{5}\partial_{\mu}\frac{\mu_{5}}{T}$$

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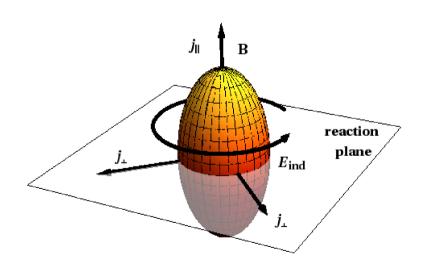
Using both hydrodynamic equations and constitutive relations one can derive

$$\partial_{\mu}(su^{\mu} - \frac{P}{T})^{\mu} - \frac{P}{T}(\nu_{5}^{\mu}) = -\frac{1}{T}(\partial_{\mu}\nu_{5})\tau^{\mu\nu} - \nu^{\mu}(\partial_{\mu}\frac{\mu}{T})^{\mu} + \frac{1}{T}E_{\mu}) - \nu_{5}^{\mu}(\partial_{\mu}\frac{\mu}{T})^{\mu}$$

- Entropy production is always non-negative
- No additional anomalous first-order corrections to the currents
- Only the "normal" component contributes to the entropy current, while the "superfluid" component has zero entropy

An additional electric current induced by the θ -field:

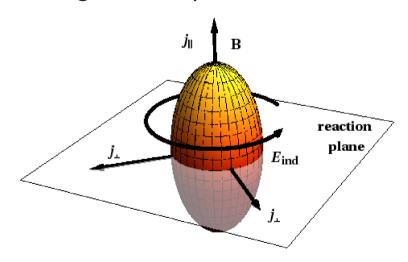
$$j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial \theta \cdot B)$$



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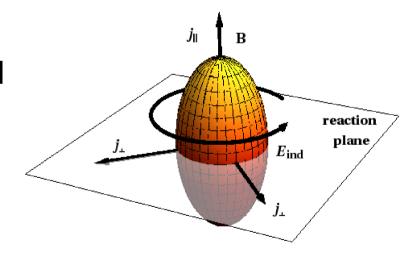
Chiral Magnetic Effect (electric current along B-field)



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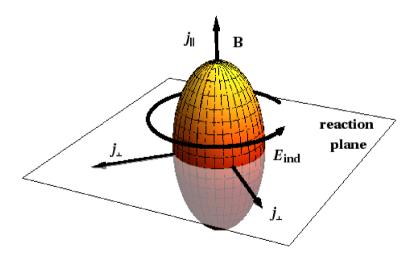
- Chiral Magnetic Effect (electric current along B-field)
- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)



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- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)
- Chiral Dipole Wave (dipole moment induced by B-field)



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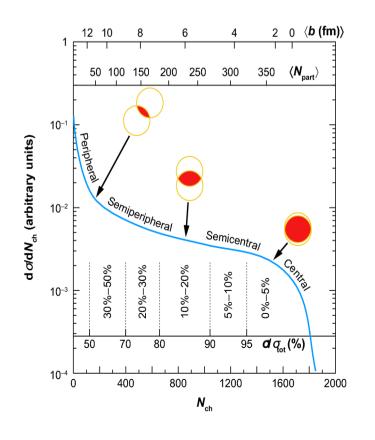
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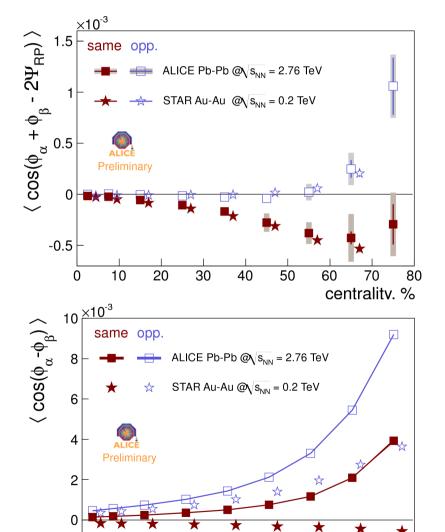
reaction plane

- Chiral Magnetic Effect (electric current along B-field)
- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)
- Chiral Dipole Wave (dipole moment induced by B-field)
- The field θ(x) itself: Chiral Magnetic Wave
 (propagating imbalance between the number of left- and right-handed quarks)

$$\gamma_{\alpha,\beta} = \langle \cos(\phi_{\alpha} + \phi_{\beta}) \rangle = \langle \cos \cos \rangle - \langle \sin \sin \rangle$$

$$\delta_{\alpha,\beta} = \langle \cos(\phi_{\alpha} - \phi_{\beta}) \rangle = \langle \cos \cos \rangle + \langle \sin \sin \rangle$$

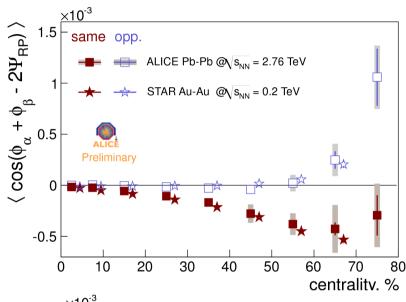


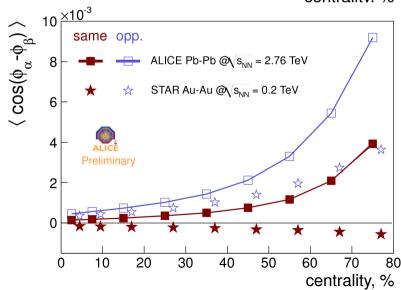


centrality, %

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in-plane out-of-plane



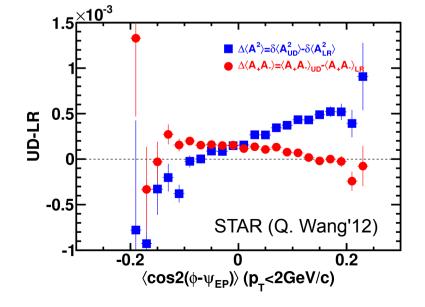


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in-plane out-of-plane

$$\gamma_{\alpha,\beta} \sim v_2 F_{\alpha,\beta} - H_{\alpha,\beta}^{\text{out}} + H_{\alpha,\beta}^{\text{in}}$$

$$\delta_{\alpha,\beta} \sim F_{\alpha,\beta} + H_{\alpha,\beta}^{\text{out}} + H_{\alpha,\beta}^{\text{in}}$$



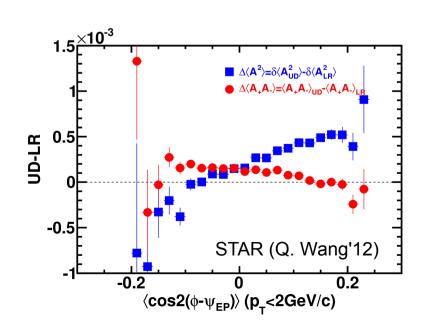
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flow-dependent flow-independent



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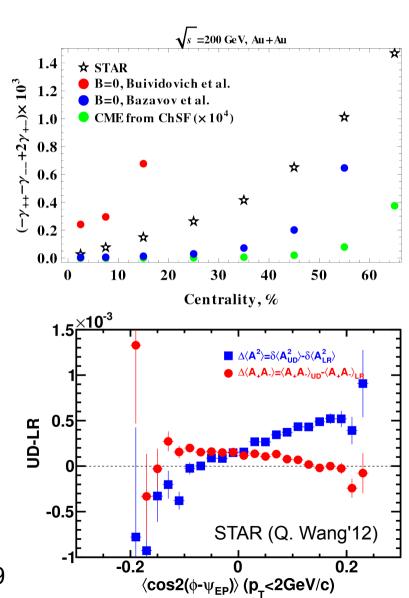
$$\gamma_{lpha,eta} \sim v_2 \, F_{lpha,eta} - H_{lpha,eta}^{
m out} + H_{lpha,eta}^{
m in}$$

$$\delta_{lpha,eta} \sim F_{lpha,eta} + H_{lpha,eta}^{
m out} + H_{lpha,eta}^{
m in}$$

flow-dependent flow-independent

$$H_{++} + H_{--} - 2H_{+-} \sim \frac{4\pi\tau^2\rho^2\mathcal{R}^2}{3N_q^2} \left(\langle j_{\parallel}^2 \rangle + 2\langle j_{\perp}^2 \rangle \right)$$

Buividovich, Chernodub, Luschevskaya, Polikarpov' 09



Interesting projects

- Add more flavors. The "axion-like" field might be just a pion. At strong B it will be then a 2D pion.
- Role of the low-dimensional defects in inducing superfluidity/superconductivity. Copenhagen vacuum.
- Entropy and thermalization of various parts of fermionic spectrum. Relation to the vacuum entropy.
- Considering vortices (axionic strings). One might reproduce the chiral vortical effect.
- Holographic model of the chiral superfluidity and highorder corrections.
- Experimental searches for the mentioned effects.
- Chiral electric effect on a lattice.

Thank you for the attention!

and

Have a good time!

All comments on the papers are welcome!
Also feel free to ask questions about the experimental observables.

Backup slide

