

# QCD string interactions and implications for high energy collisions

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**1503.05213:** Collective flow in high-multiplicity proton-proton collisions

**1404.1888:** Collective interaction of QCD strings and early stages of high multiplicity pA collisions

**1402.7363:** Self-interacting QCD strings and string balls



# What systems do we consider?

- **AA collision and QCD in the mixed phase**
  - **high multiplicity pA collisions**
  - **peripheral AA collisions**
  - **high multiplicity pp collisions**
  - **ultra-high energy cosmic rays**
- (if time permits)**

# Spaghetti

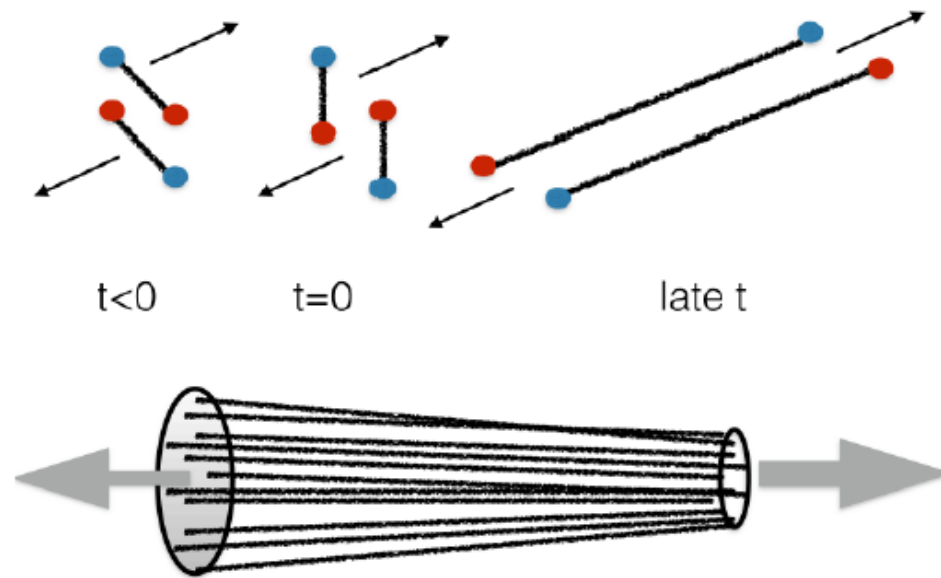


FIG. 1: The upper plot reminds the basic mechanism of two string production, resulting from color reconnection. The lower plot is a sketch of the simplest multi-string state, produced in  $pA$  collisions or very peripheral  $AA$  collisions, known as “spaghetti”.

# Motivation: multiplicity in pA

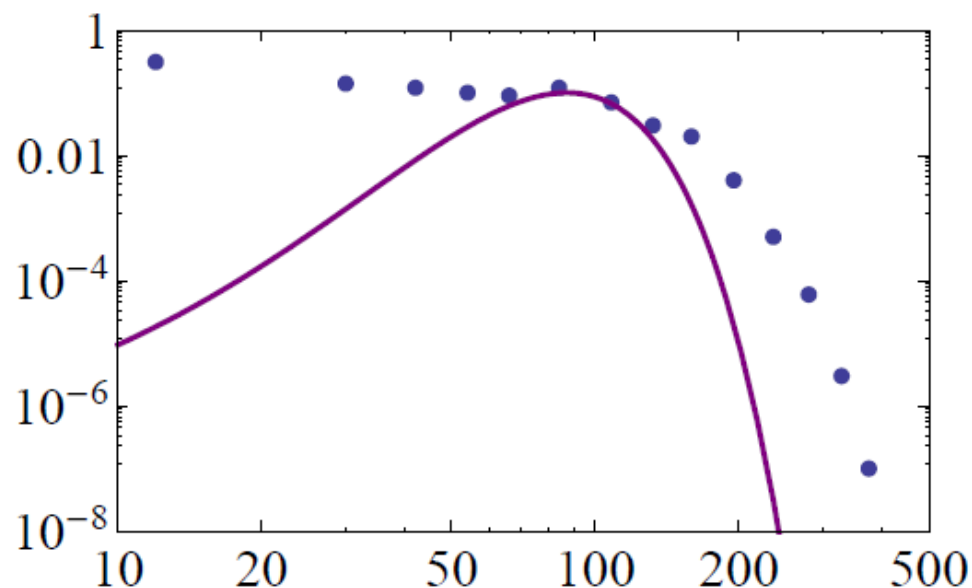
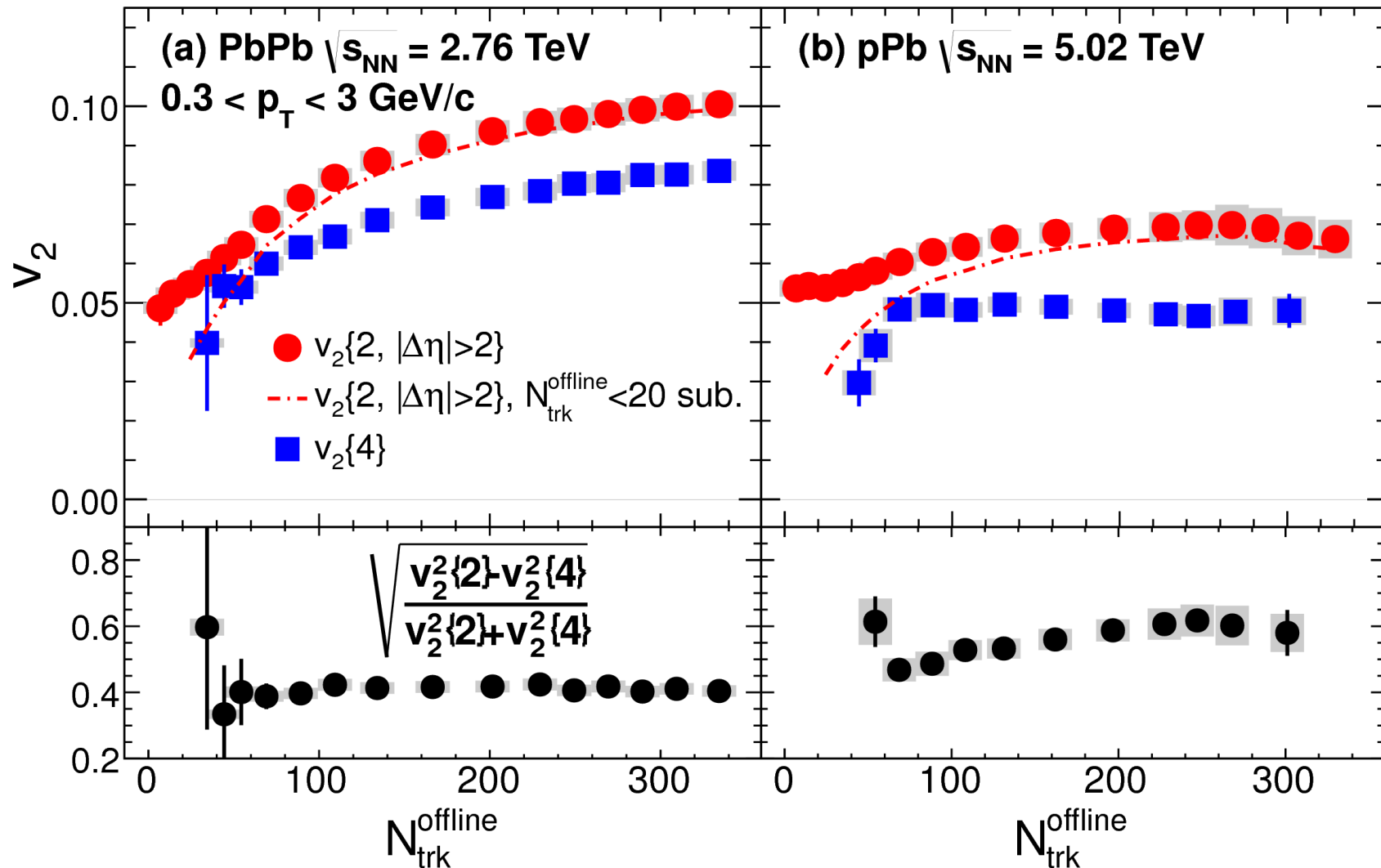
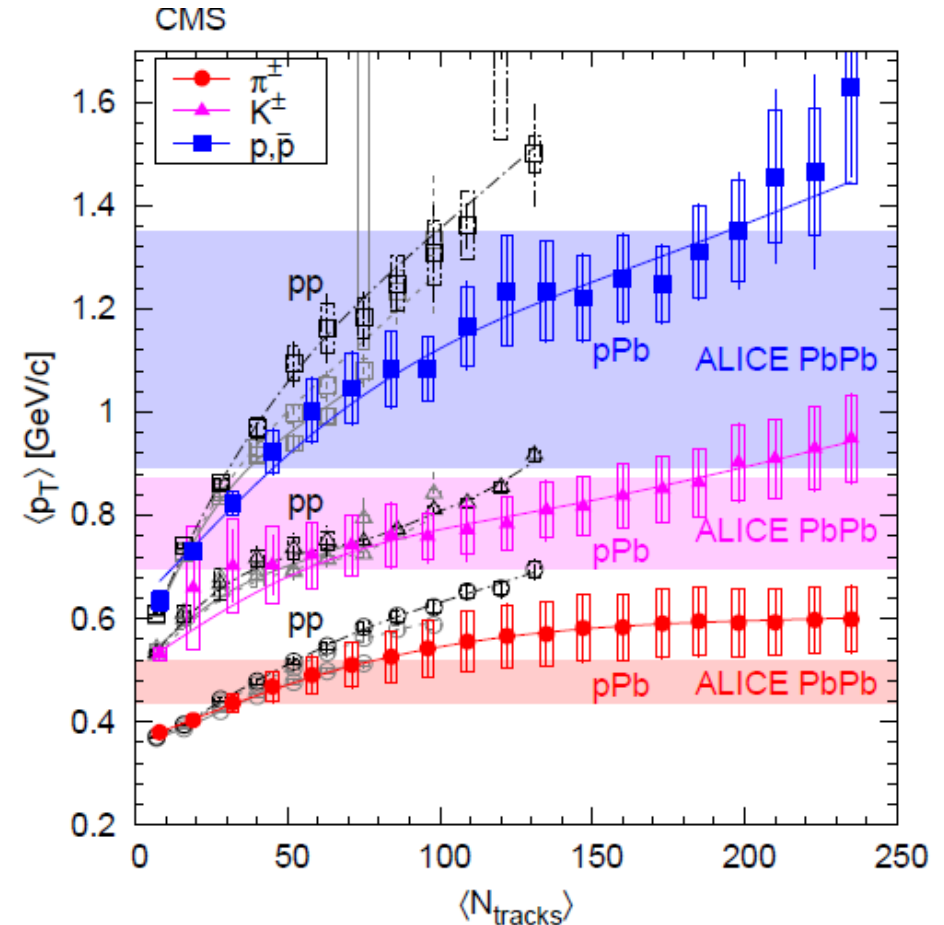
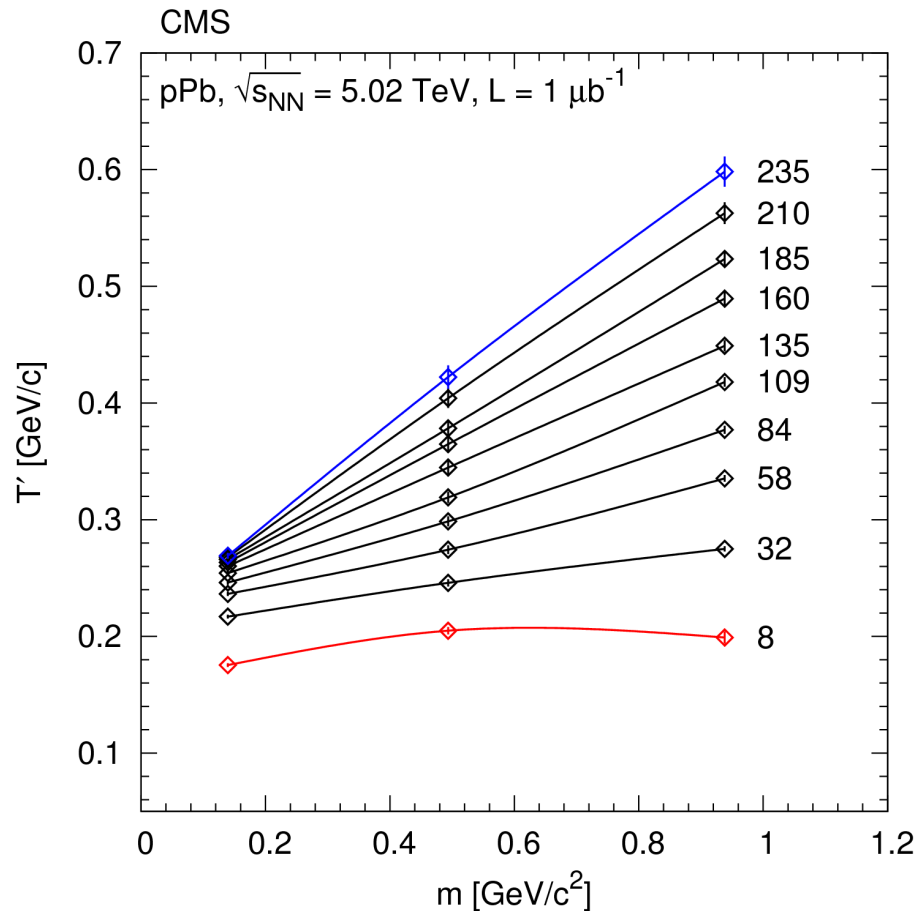


FIG. 2: (Color online). Probability distribution over the number of charged tracks in the CMS detector acceptance  $P(N_{tr})$  [13]. The (purple) line is the Poisson distribution with  $\langle N_p \rangle = 16$ , arbitrarily normalized to touch the data points.

# Motivation: elliptic flow in pPb



# Motivation: radial flow in pPb



$$\frac{dN}{dy dp_\perp^2} \sim \exp \left( - \frac{\sqrt{m^2 + p_\perp^2}}{T'} \right)$$

**Inverse slope parameters  $T'$  from fits of pion, kaon, and proton spectra (both charges)**

# History: collapse to a black hole

Susskind; Horowitz and Polchinski; Damour and Veneziano.

**Free string:**

$$S_{string} \sim M/M_s$$

$$\frac{R_{ball, r.w.}}{l_s} \sim \sqrt{M}$$

**S(string)=S(BH) only at some special mass.**

**Black hole:**

**Temperature  $T=T_H$ .**

**But the sizes don't match, so we need self-interaction.**

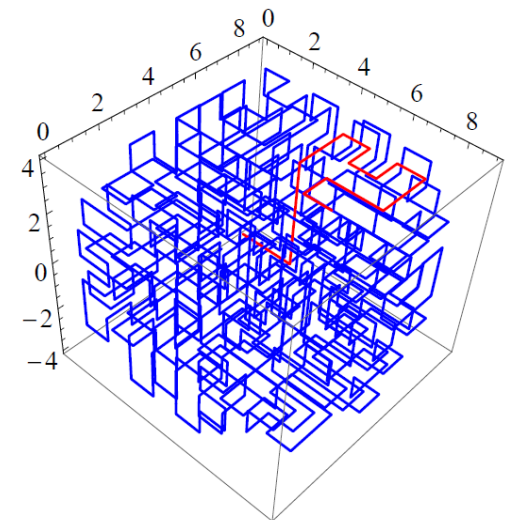
$$S_{BH} \sim Area \sim M^{\frac{d-1}{d-2}}$$

$$R_{BH} \sim (M)^{\frac{1}{(d-2)}}$$

**String ball:**

$$S(M, R) \sim M \left(1 - \frac{1}{R^2}\right) \left(1 - \frac{R^2}{M^2}\right) \left(1 + \frac{g^2 M}{R^{d-2}}\right)$$

**Something like this in QCD?**



# Hagedorn phenomenon

**Partition function for strings on a lattice**

$$Z \sim \int dL \exp \left[ \frac{L}{a} \ln(2d - 1) - \frac{\sigma_T L}{T} \right]$$



# Hagedorn phenomenon

Partition function for strings on a lattice

$$Z \sim \int dL \exp \left[ \frac{L}{a} \ln(2d-1) - \frac{\sigma_T L}{T} \right]$$

Energy

Entropy factor

Hagedorn transition temperature (zero effective tension of the string)

$$T_H = \frac{\sigma_T a}{\ln(2d-1)}$$

Bringoltz & Teper '06:  $T_H/T_c = 1.11$

What happens with the string at the critical temperature? Let's put in on a lattice.

$$a \simeq 0.54 \text{ fm} \quad E_{pl} = 4\sigma_T a \simeq 1.9 \text{ GeV} \quad \frac{\epsilon_{max}}{T_c^4} = \frac{\sigma_T a}{a^3 T_c^4} \approx 4.4$$

$$\sigma_T = (0.42 \text{ GeV})^2 \quad E_m = \sigma_T a \simeq 0.5 \text{ GeV} \quad \frac{\epsilon_{gluons}}{T^4} = (N_c^2 - 1) \frac{\pi^2}{15} \approx 5.26$$

# String on a lattice

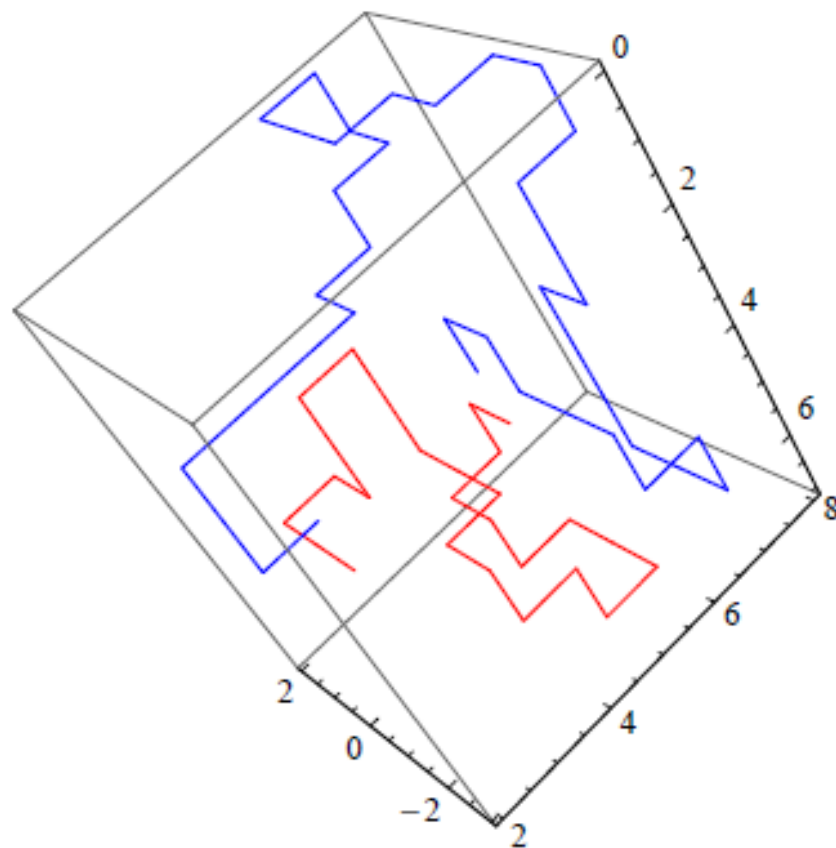


FIG. 4: (Color online) Example of a two-string configuration (a sparse string ball): two strings are plotted as blue and red.

# Sigma-cloud

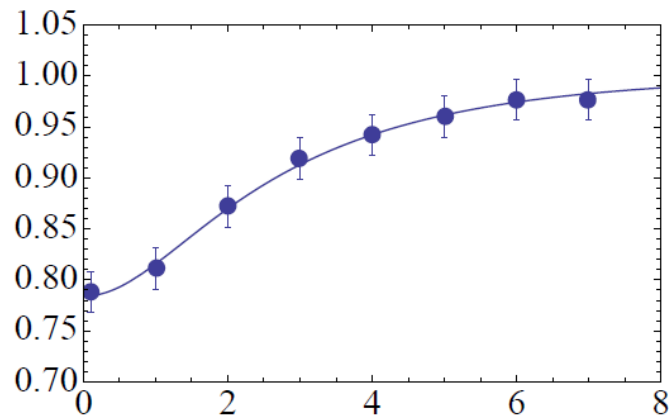
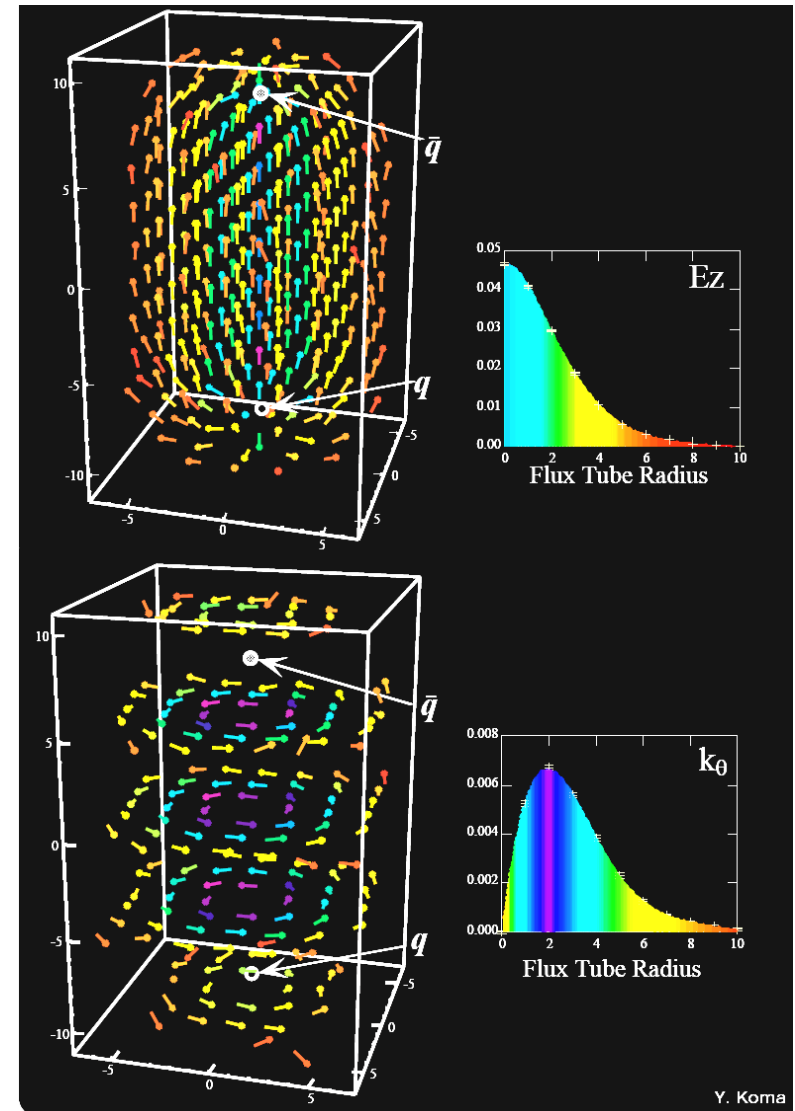


FIG. 3: (Color online). Points are from the lattice data for the chiral condensate [16]. The curve is expression (7) with  $C = 0.26$ ,  $s_{string} = 0.176$  fm.

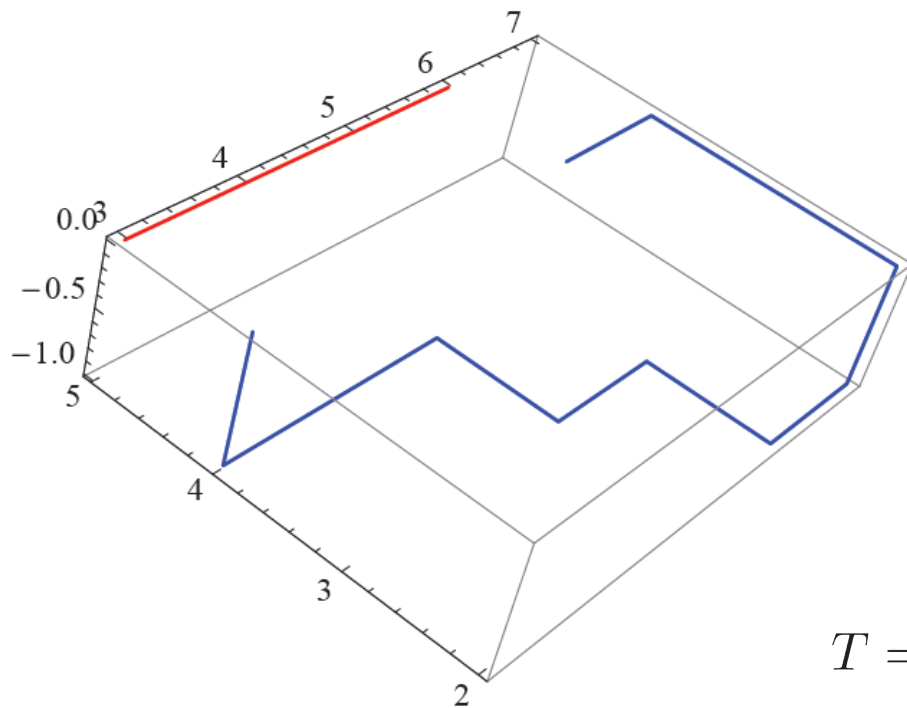
$$\frac{\langle \bar{q}q(r_{\perp})W \rangle}{\langle W \rangle \langle \bar{q}q \rangle} = 1 - CK_0(m_{\sigma}\tilde{r}_{\perp}),$$

$$\tilde{r}_{\perp} = \sqrt{r_{\perp}^2 + s_{string}^2}$$

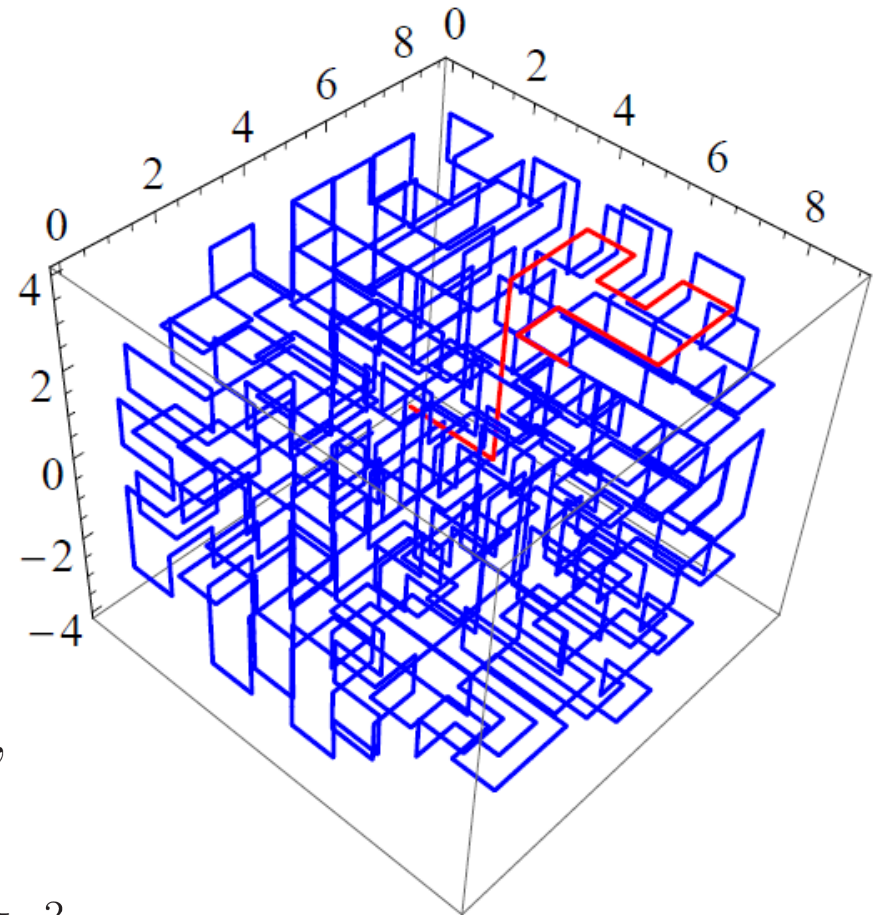
**Type I dual superconductor**



# Interacting strings



**Without self-interaction**



**With self-interaction**

$$\begin{aligned}T &= 1 \text{ GeV}, \\s_T &= 1.5a, \\g_N &= 4.4 \text{ GeV}^{-2}\end{aligned}$$

# String balls

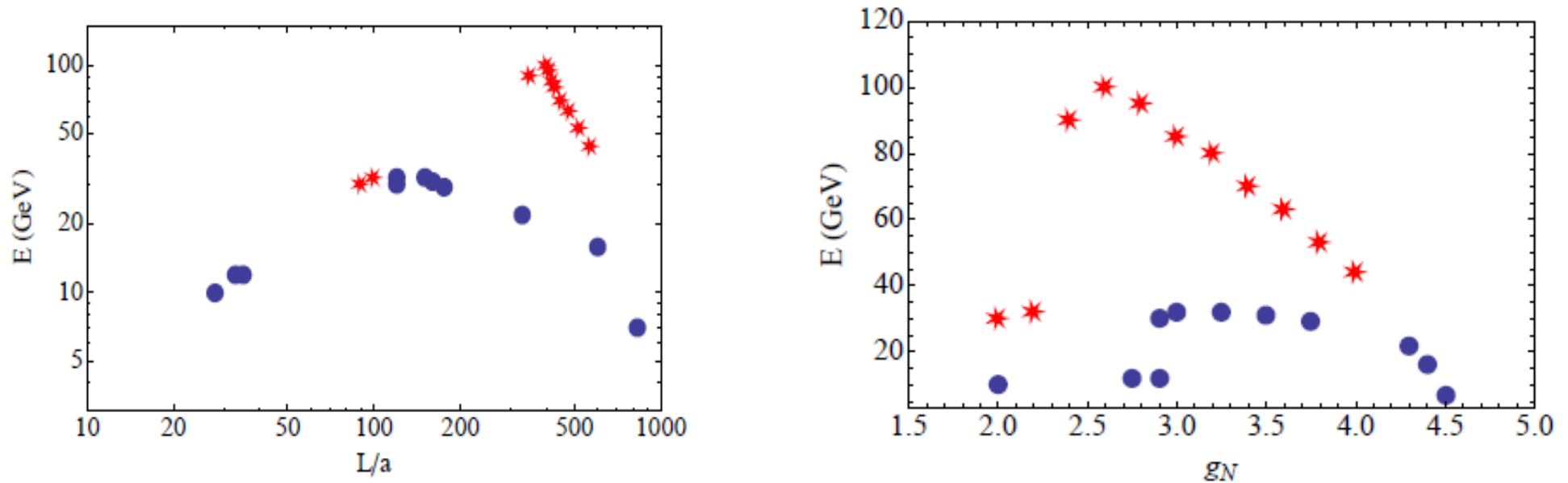


FIG. 7: Upper plot: The energy of the cluster  $E$  (GeV) versus the length of the string  $L/a$ . Lower plot: The energy of the cluster  $E$  (GeV) versus the “Newton coupling”  $g_N$  ( $\text{GeV}^{-2}$ ). Points show the results of the simulations in setting  $T_0 = 1$  GeV and size of the ball  $s_T = 1.5a, 2a$ , for circles and stars, respectively.

## Applications:

1. Jet queching

2. Angular correlations

# Jet quenching

Jet quenching parameter, by definition

$$\hat{q} = \frac{d\langle p_{\perp}^2 \rangle}{dl}$$

in our case (due to „kicks“ by the color force)

$$\hat{q} \approx \frac{16}{3} \alpha_s \sigma_T \frac{\bar{L} r_s}{\text{fm}^3}$$

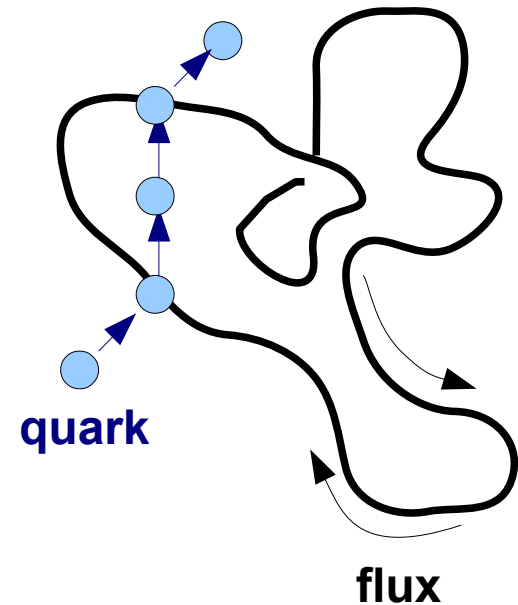
In numbers, in the mixed phase

(min-max are because of the string density per cubic fermi):

$$\hat{q}_{min} = 0.028, \quad \hat{q}_{max} = 0.10 \left( \frac{\text{GeV}^2}{\text{fm}} \right)$$

Compare to the data of JET collaboration (1312.5003):

$$\hat{q}_{min} = 0.025, \quad \hat{q}_{max} = 0.15 \left( \frac{\text{GeV}^2}{\text{fm}} \right)$$



**Can be up to 1 GeV<sup>2</sup>/fm  
for a string ball!**

# Angular correlations

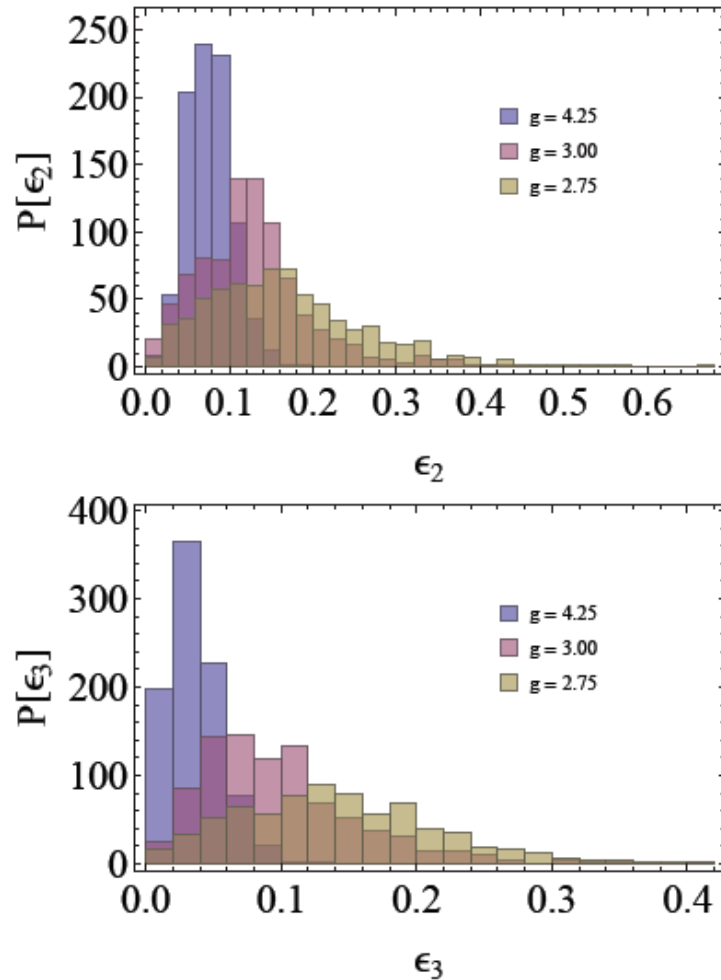


FIG. 10: The distributions over  $\epsilon_2$  and  $\epsilon_3$  (upper and lower plots), for several values of the “Newton coupling”  $g_N$  [GeV $^{-2}$ ].

$$\epsilon_n = \frac{\int d^2 r_{\perp} \cos(n\phi) r_{\perp}^n (dN/d^2 r_{\perp})}{\int d^2 r_{\perp} r_{\perp}^n (dN/d^2 r_{\perp})}$$

$$(\epsilon_n\{2\})^2 = \langle \epsilon_n^2 \rangle,$$

$$(\epsilon_n\{4\})^4 = 2\langle \epsilon_n^2 \rangle^2 - \langle \epsilon_n^4 \rangle,$$

$$(\epsilon_n\{6\})^6 = \frac{1}{4} [\langle \epsilon_n^6 \rangle - 9\langle \epsilon_n^2 \rangle \langle \epsilon_n^4 \rangle + 12\langle \epsilon_n^2 \rangle^3],$$

$$\epsilon_2\{2\} = 0.0759, \quad \epsilon_2\{4\} = 0.0621,$$

$$\epsilon_2\{6\} = 0.0636, \quad \epsilon_2\{8\} = 0.0635$$

$$v_2\{2\} > v_2\{4\} \approx v_2\{6\} \approx v_2\{8\}$$

# Spaghetti

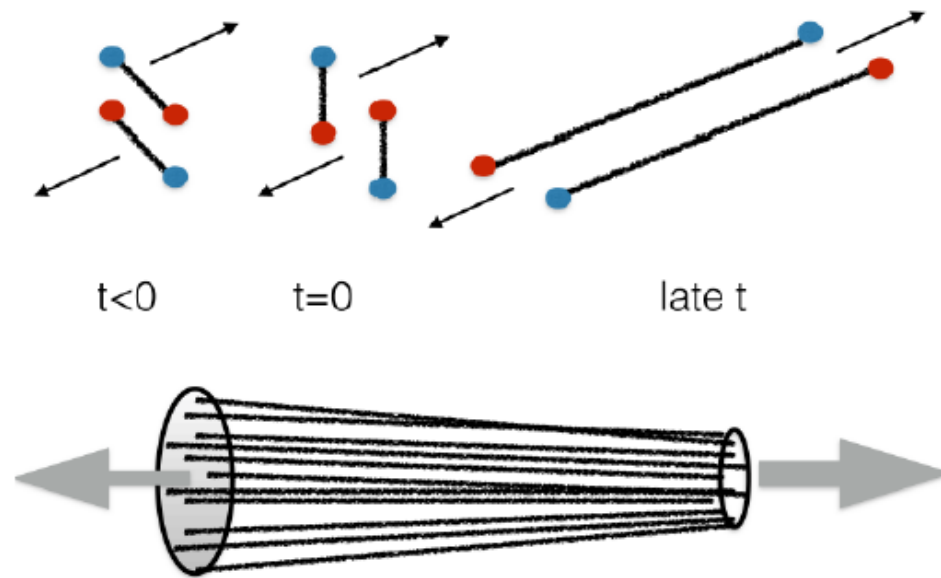


FIG. 1: The upper plot reminds the basic mechanism of two string production, resulting from color reconnection. The lower plot is a sketch of the simplest multi-string state, produced in  $pA$  collisions or very peripheral  $AA$  collisions, known as “spaghetti”.



# 2D Yukawa gas

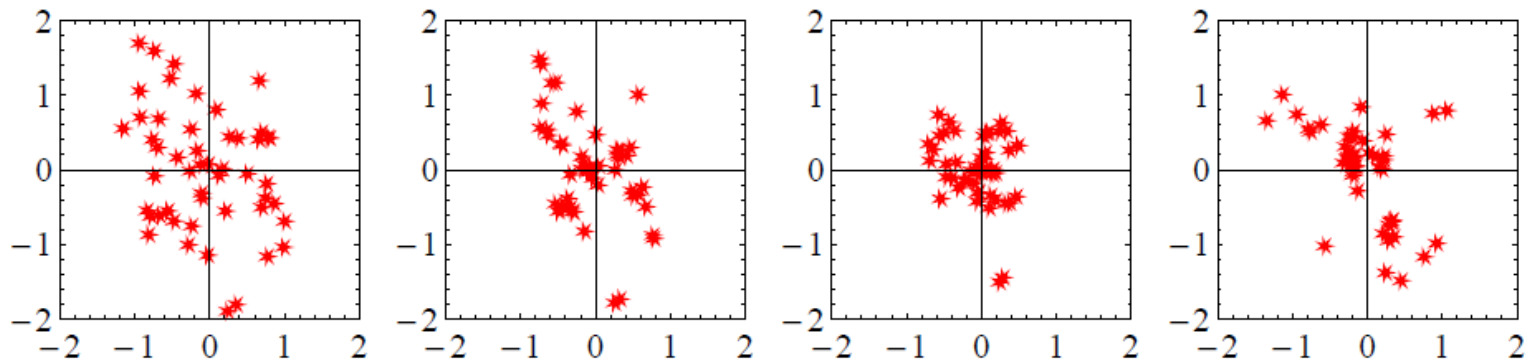


FIG. 7: (Color online) Example of changing transverse positions of the 50 string set: four pictures correspond to one initial configuration evolved to times  $\tau = 0.1, 0.5, 1, 1.5$  fm/c. The distances are given in fm, and  $g_N \sigma_T = 0.2$ .

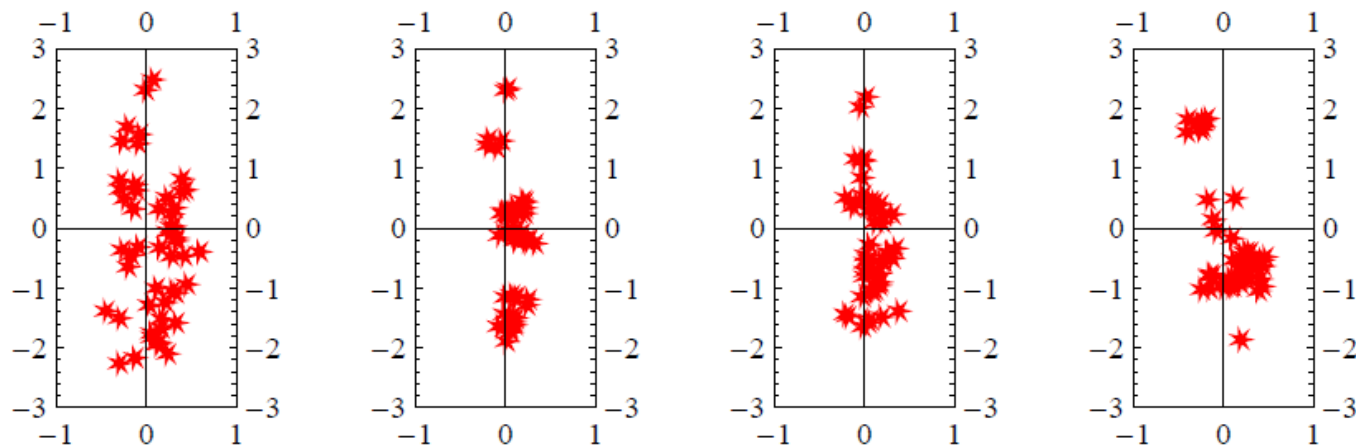


FIG. 8: (Color online) Example of peripheral AA collisions, with  $b = 11$  fm,  $g_N \sigma_T = 0.2$ , and the 50 string set. Four snapshots of the string transverse positions  $x, y$  (fm) correspond to times  $\tau = 0.1, 0.5, 1, 2.6$  fm/c.

# Energy and energy density

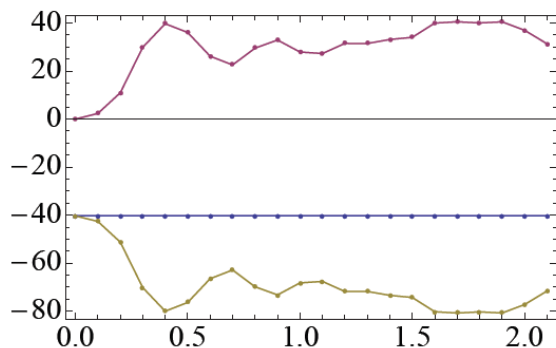


FIG. 5: (Color online). The (dimensionless) kinetic and potential energy of the system (upper and lower curves) for the same example as shown in Fig. 7, as a function of time  $t$  (fm/c). The horizontal line with dots is their sum,  $E_{tot}$ , which is conserved.

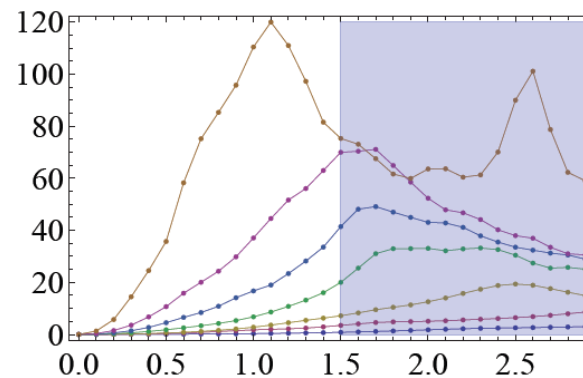


FIG. 6: (Color online). Kinetic energy (dimensionless) versus the simulation time (fm/c), for few  $pA$   $N_s = 50$  runs. Seven curves (bottom-to-top) correspond to increasing coupling constants  $g_N \sigma_T = 0.01, 0.02, 0.03, 0.05, 0.08, 0.10, 0.20$ .

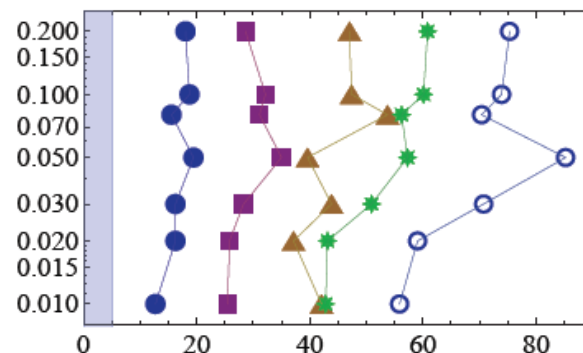
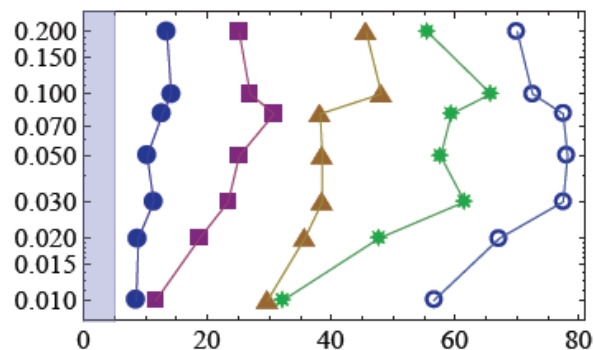


FIG. 9: (Color online) The left plot is for central  $pA$ , the right one – for peripheral  $AA$  collisions. The vertical axis is the effective coupling constant  $g_N \sigma_T$  (dimensionless). The horizontal axis is the maximal energy density  $\epsilon_{max}$  (GeV/fm<sup>3</sup>) defined by the procedure explained in the text. Five sets shown by different symbols correspond to string number  $N_s = 10, 20, 30, 40, 50$ , left to right respectively.

# Elliptic flow

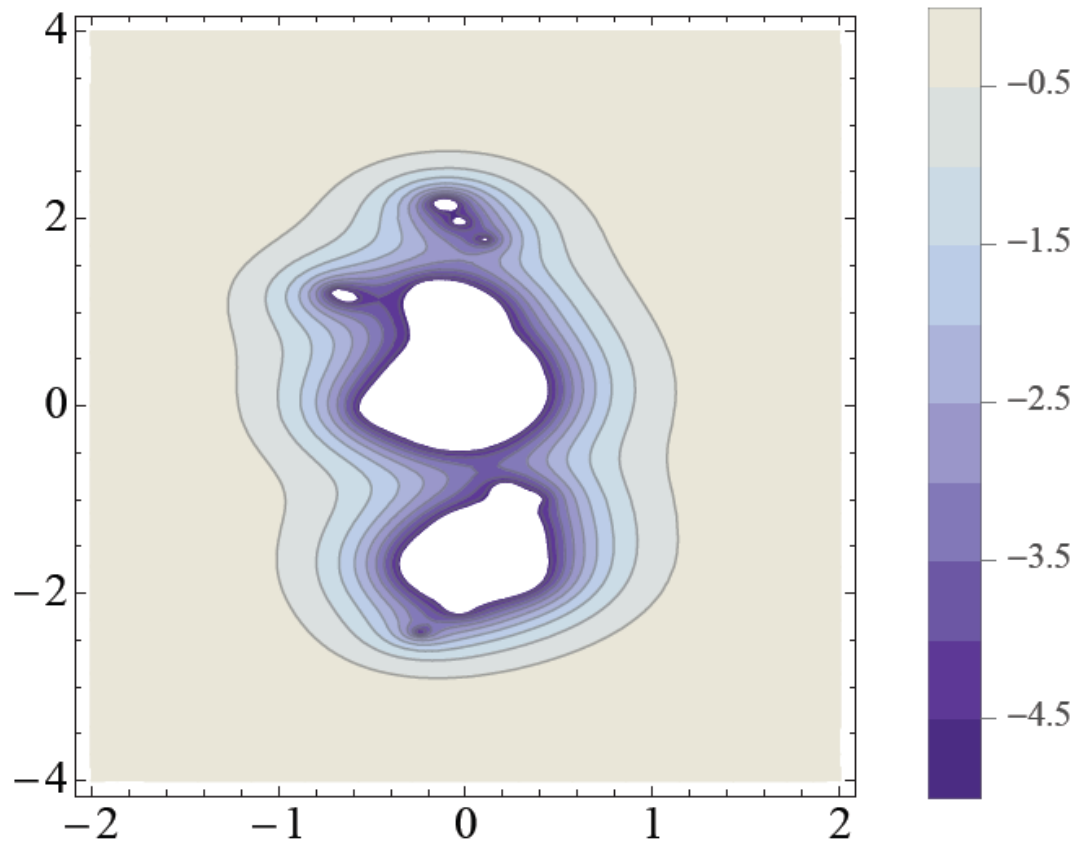
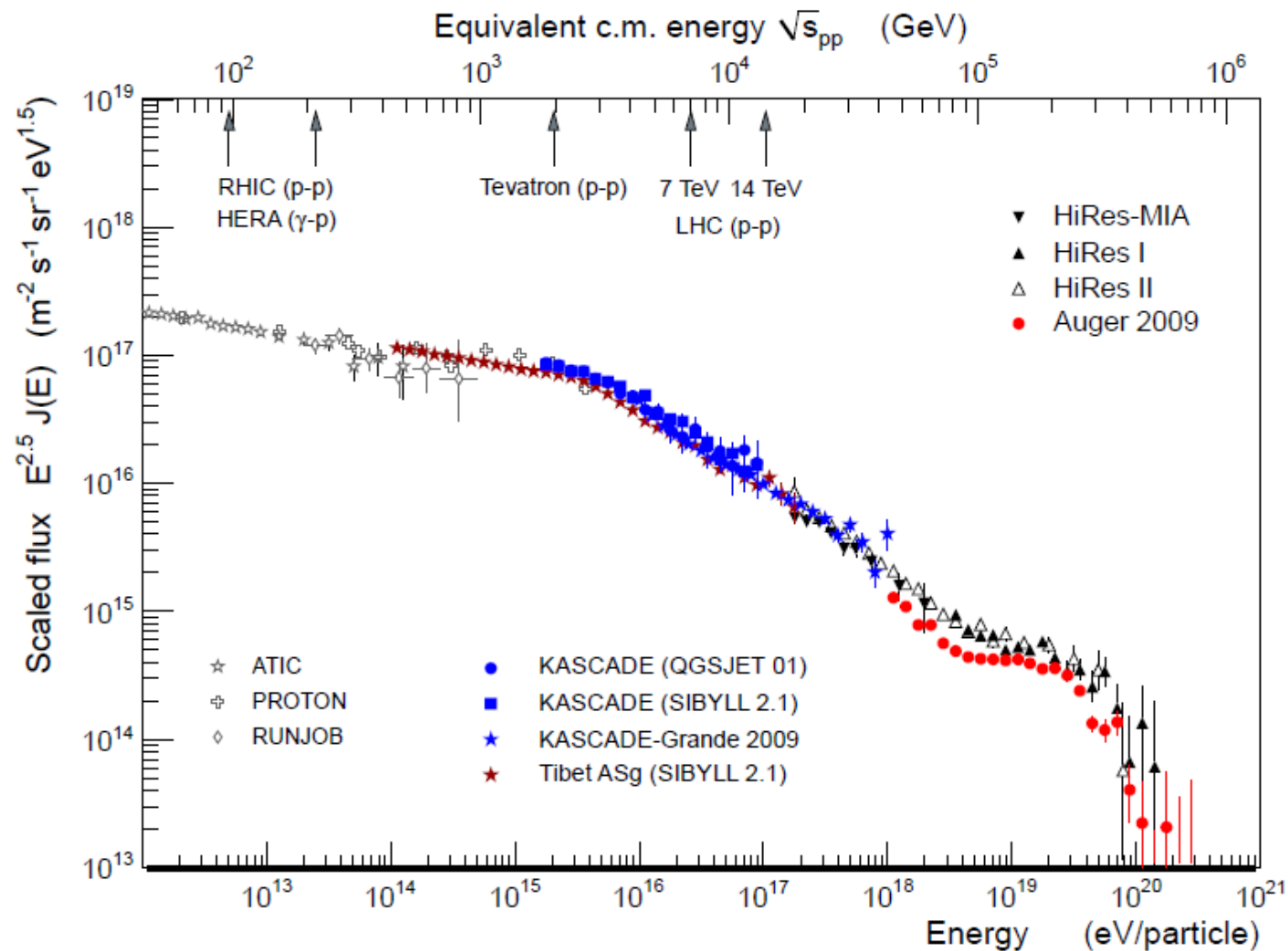
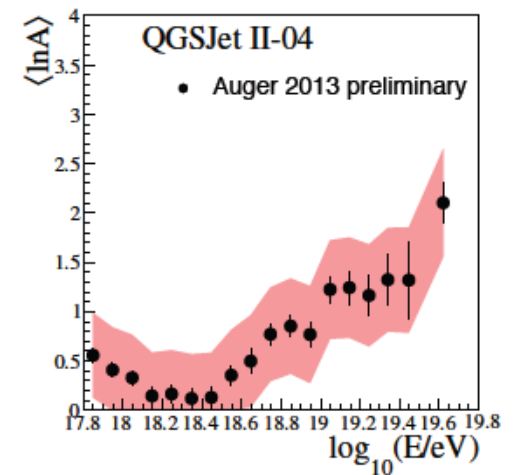


FIG. 10: Instantaneous collective potential in units  $2g_N\sigma_T$  for an  $AA$  configuration with  $b = 11$  fm,  $g_N\sigma_T = 0.2$ ,  $N_s = 50$  at the moment of time  $\tau = 1$  fm/ $c$ . White regions correspond to the chirally restored phase.

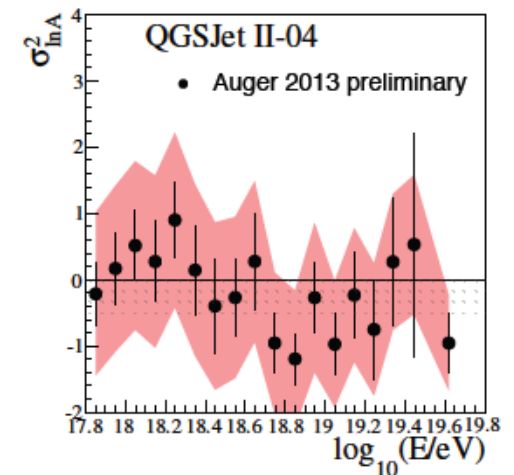
# Cosmic rays



1101.5596



Fe



p

1310.4620

# Freezeout surfaces

p O

Fe O

Pb Pb

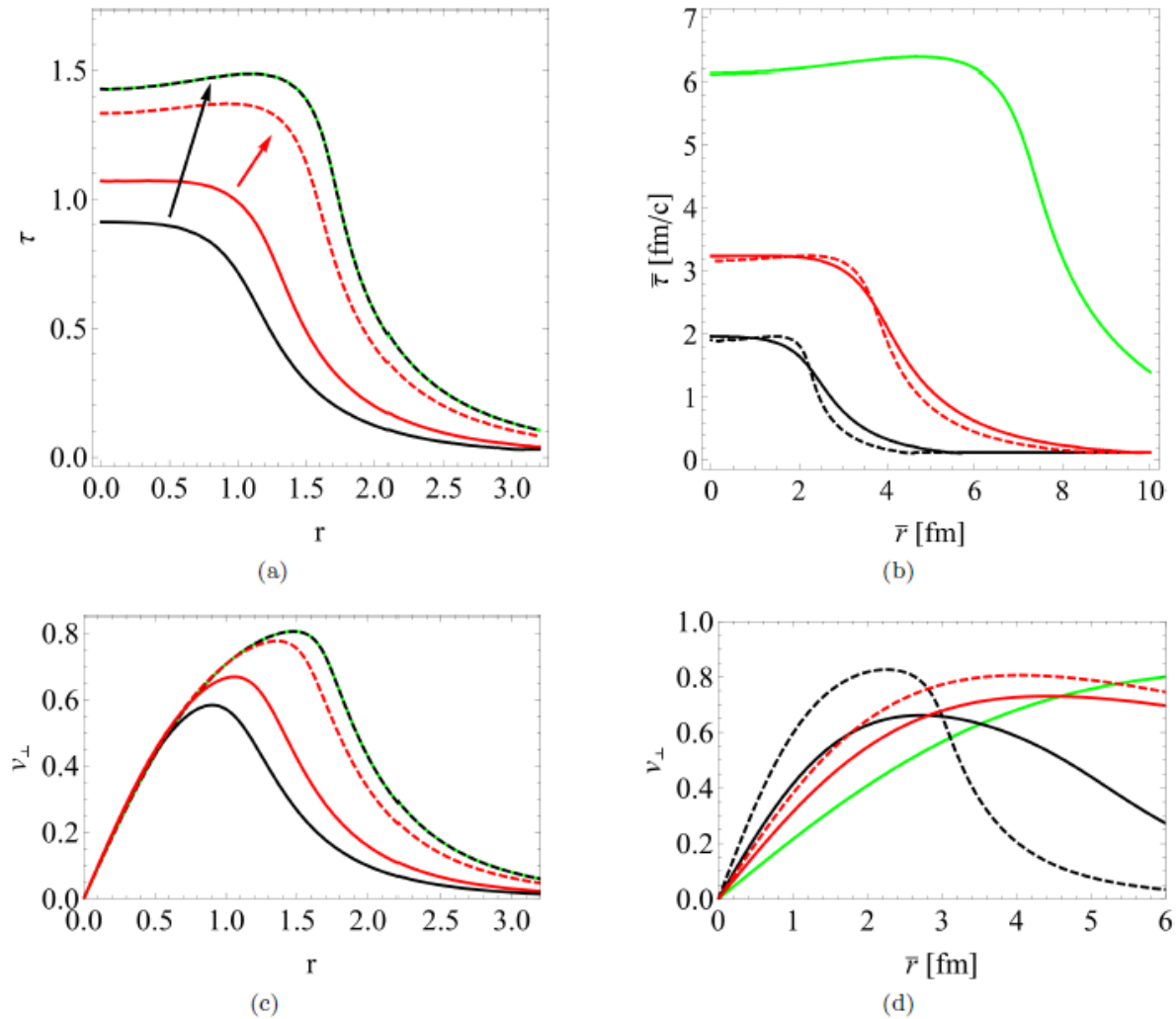
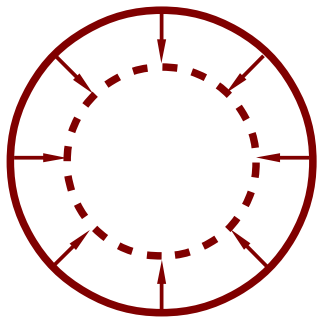


FIG. 2: (color online) (a) The freezeout surfaces in the  $(\tau, r)$  plane and (c) the distribution of the transverse flow velocity on those surfaces. In both plots the green solid curve at the top is our “benchmark”, the central *PbPb* collisions at LHC. Black solid line is for light-light collisions, black dashed (coincident with green by chance) are light-light collisions with the size compression. Similarly, red solid and red dashed are heavy-light collisions without and with the size compression, respectively.

# Particle spectra

From Cooper-Frye formula and freezeout curves (1407.3270)

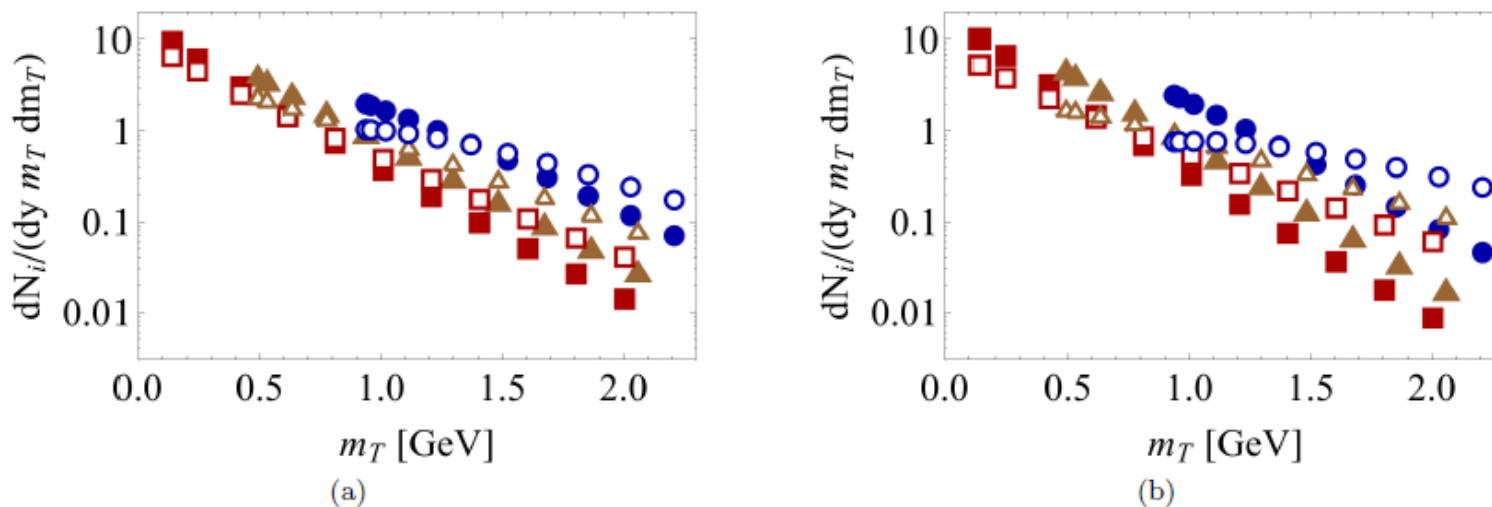


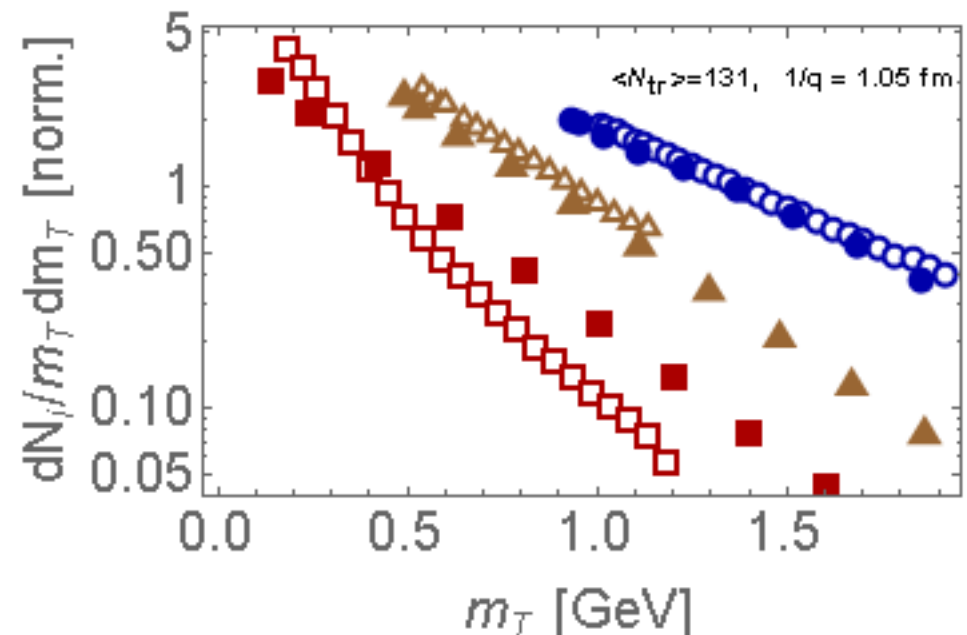
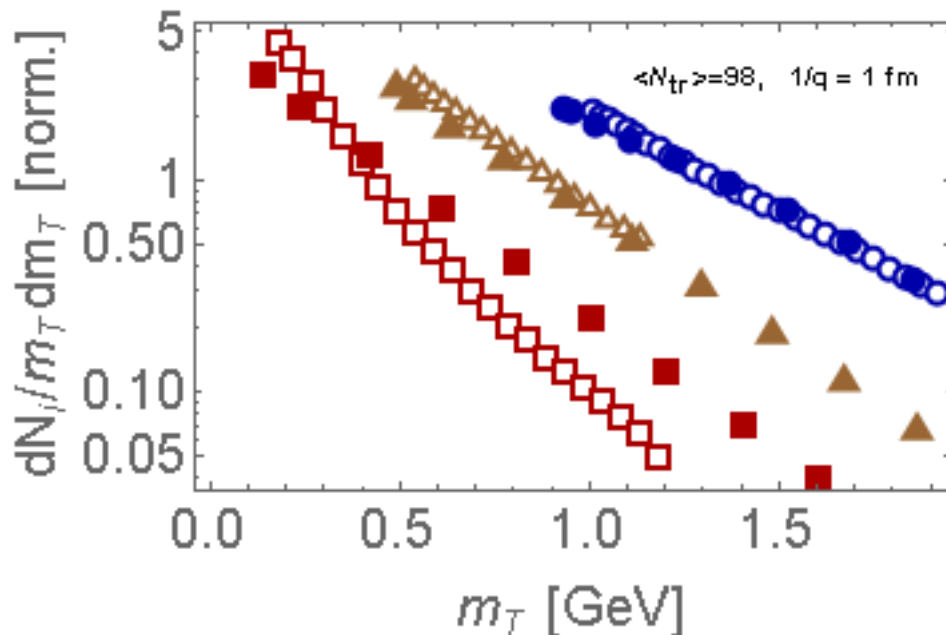
FIG. 3: (color online) Normalized spectra of pions (squares), kaons (triangles) and protons (discs) for the (a) heavy-light (e.g. *FeO*) and (b) light-light (e.g. *pO*) collisions. Open symbols correspond to the “compressed” cases, explained in the text.

**Mean  $p_T$  :**

particles	FeO	FeO comp.	pO	pO comp.	PbPb
$\pi^\pm$	0.56	0.69	0.53	0.76	0.73
$K^\pm$	0.71	0.88	0.66	0.96	0.92
$p, \bar{p}$	0.90	1.09	0.83	1.17	1.13

# pp with Gubser's flow

From Cooper-Frye formula and freezeout curves



The data (open symbols) are from CMS, 7 TeV,  
fit done by the system size parameter  $q$  in the Gubser's flow solution.  
**protons**, **kaons**, **pions**

How to make a transition to QGP/hydro? String balls?

# Conclusions

- **One should reconsider the QCD string phenomenology taking into account the interaction between strings. More lattice data are needed.**
- **Jet quenching in the inhomogeneous phases.**
- **One should implement the interaction in order to describe the collective effects in pA collisions. The Lund model based approaches may be improved.**
- **Naive energy extrapolation of the LHC results (and Monte-Carlo generators) for ultra-high-energy cosmic rays should be corrected.**
- **There might be hydro in pp-collisions.**



**Thank you for the  
attention!**