

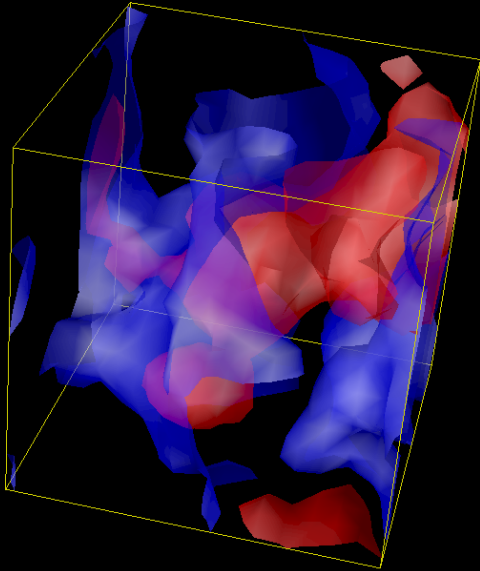
Fluid/gravity model for the Chiral Magnetic Effect



Tigran Kalaydzhyan

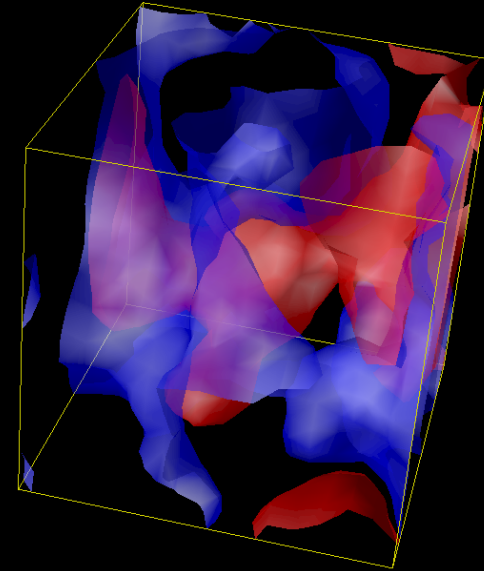
Spring School on Superstring Theory and Related Topics

28 March - 05 April 2011. The Abdus Salam International Centre for
Theoretical Physics, Trieste, Italy.

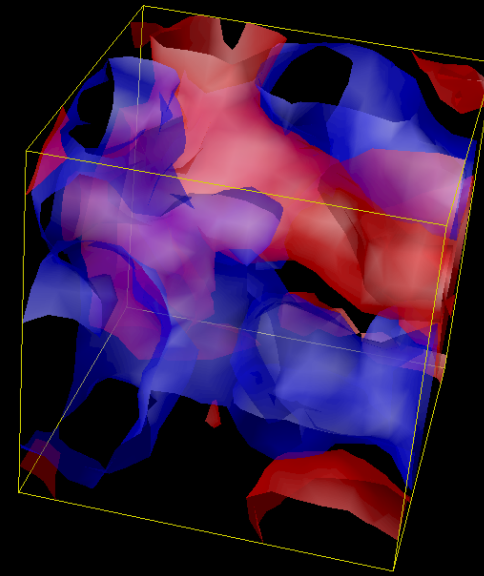
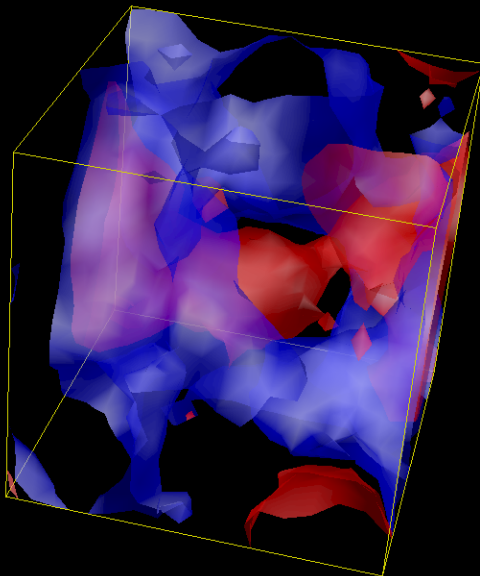


Negative topological
charge density

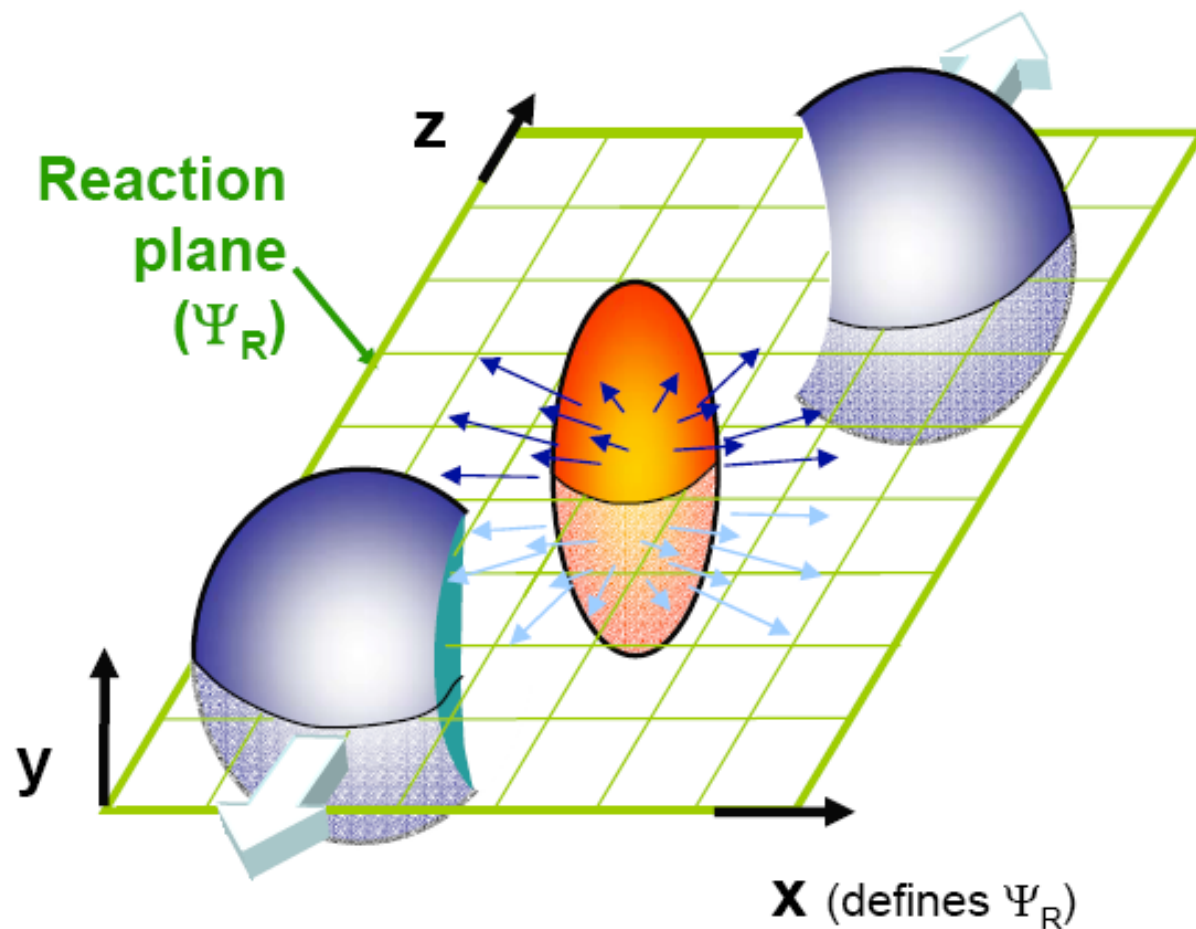
QCD Vacuum



Positive topological
charge density



Chiral Magnetic Effect

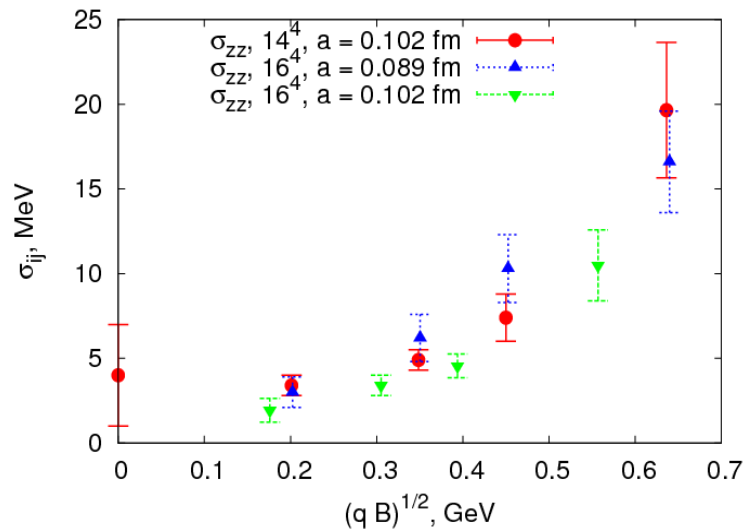
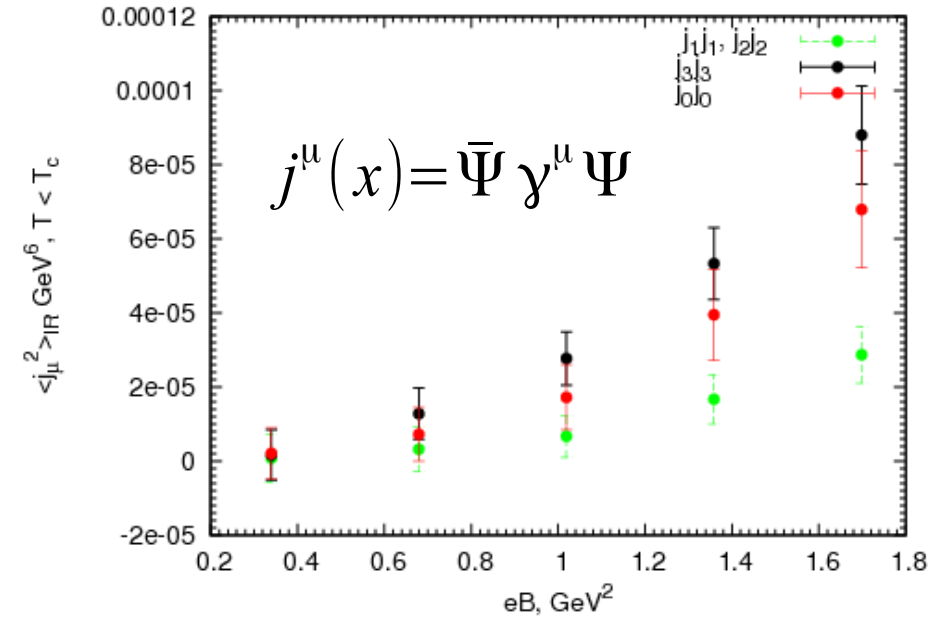
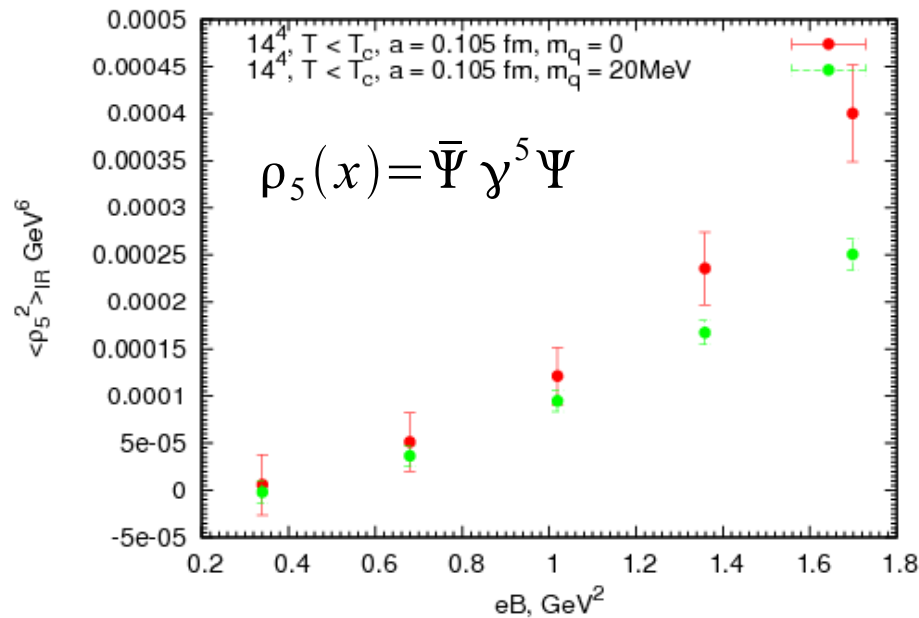


Excess of positive charge

$$J^\mu = \kappa_B B^\mu$$

Excess of negative charge

Some numbers (lattice)

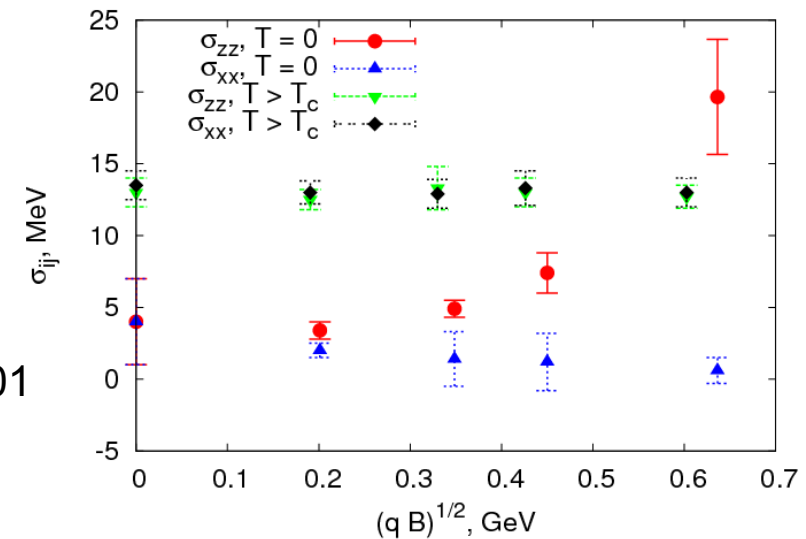


T.K., D. Kharzeev and

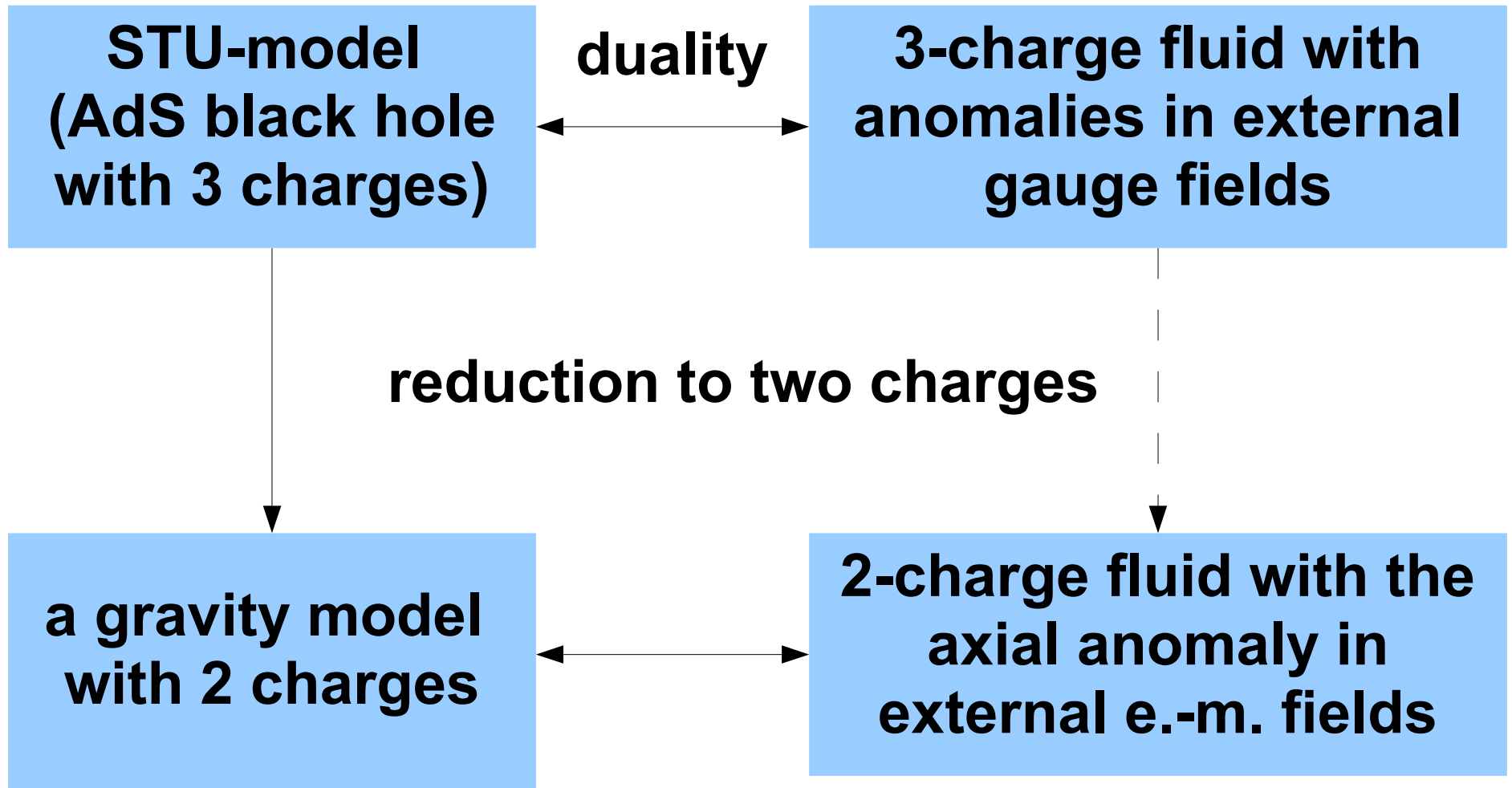
ITEP → Lattice

PRL 105 (2010) 132001

PoS LAT (2010) 190



Main idea



Hydrodynamics

Three-charge model:

$$\partial_\mu T^{\mu\nu} = F^{a\nu\lambda} j_\lambda^a, \quad a=1,2,3$$

$$\partial_\mu j^{a\mu} = -\frac{1}{8} C^{abc} F_{\mu\nu}^b \tilde{F}^{c\mu\nu} = C^{abc} E^b \cdot B^c$$

where stress-energy tensor and U(1) currents:

$$T^{\mu\nu} = (\epsilon + P) u^\mu u^\nu + P g^{\mu\nu} + \dots,$$

$$j^{a\mu} = \rho^a u^\mu + \xi_\omega^a \omega^\mu + \xi_B^{ab} B^{b\mu} + \dots$$

Electric field

$$E^{a\mu} = u_\nu F^{a\mu\nu}$$

Magnetic field

$$B^{a\mu} = \frac{1}{2} \epsilon^{\mu\nu\lambda\rho} u_\nu F_{\lambda\rho}^a$$

Vorticity

$$\omega^\mu = \frac{1}{2} \epsilon^{\mu\nu\lambda\rho} u_\nu \partial_\lambda u_\rho$$

Quantum anomaly → classical dynamics!

Son and Surowka (2009)

Transport coefficients

$$j^{a\mu} = \rho^a u^\mu + \xi_\omega^a \omega^\mu + \xi_B^{ab} B^{b\mu} + \dots$$

where the coefficients are

$$\xi_\omega^a = C^{abc} \mu^b \mu^c - \frac{2}{3} \rho^a C^{bcd} \frac{\mu^b \mu^c \mu^d}{\epsilon + P}$$

$$\xi_B^{ab} = C^{abc} \mu^c - \frac{1}{2} \rho^a C^{bcd} \frac{\mu^c \mu^d}{\epsilon + P}$$

Here μ^a is a chemical potential associated with density ρ^a

Reduction to two charges

Hydrodynamic equations:

$$\partial_\mu T^{\mu\nu} = F^{\nu\lambda} j_\lambda,$$

$$\partial_\mu j_5^\mu = -\frac{1}{8} C F_{\mu\nu} \tilde{F}^{\mu\nu} = C E^\lambda \cdot B_\lambda,$$

$$\partial_\mu j^\mu = 0$$

where vector and axial currents are

identifications:

$$j^\mu = j^{2\mu} + j^{3\mu}$$

$$j_5^\mu = j^{1\mu}$$

$$j^\mu = \rho u^\mu + \kappa_\omega \omega^\mu + \kappa_B B^\mu + \dots$$

$$j_5^\mu = \rho_5 u^\mu + \xi_\omega \omega^\mu + \xi_B B^\mu + \dots$$

CVE $\kappa_\omega = 2C\mu\mu_5 \left(1 - \frac{\mu\rho}{\epsilon + P}\right),$

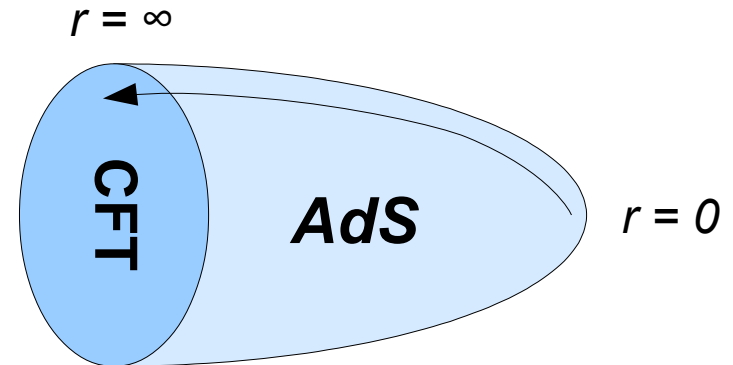
QVE $\xi_\omega = C\mu^2 \left(1 - 2\frac{\mu_5\rho_5}{\epsilon + P}\right),$

$\kappa_B = C\mu_5 \left(1 - \frac{\mu\rho}{\epsilon + P}\right),$ **CME**

$\xi_B = C\mu \left(1 - \frac{\mu_5\rho_5}{\epsilon + P}\right),$ **QME**

Holography. Algorithm.

Fluid on the boundary, gravity in the bulk. Input: energy density, anomaly, background fields, etc.



Fix metric components (and gauge fields components), Chern-Simons parameters, etc in the bulk.

Solve equations of motion for the bulk fields (i.e. Einstein-Maxwell-...).

Read off a nontrivial result (i.e. transport coefficients) from the near-boundary expansion of the gravity solution.

see also Bhattacharyya, Hubeny, Minwalla and Rangamani (2008), Torabian and Yee (2009)

Gravity. STU-model.

Holographic dual of $U(1)^3$ theory – the STU-model:

$$\mathcal{L} = R - \frac{1}{2} G_{ab} F_{MN}^a F^{bMN} - G_{ab} \partial_M X^a \partial^M X^b + \frac{1}{24 \sqrt{-g_5}} \epsilon^{MNPQR} S_{abc} F_{MN}^a F_{PQ}^b A_R^c + 4 \sum_{a=1}^3 \frac{1}{X^a}.$$

Here we have:

1. Metric g_{MN} , where $M, N = 0, \dots, 4$.

2. Three $U(1)$ gauge fields A_M^a , where $a = 1, 2, 3$.

$$G_{ab} = \frac{1}{2} \delta_{abc} (X^c)^{-2}$$

3. Three scalars X^a : $X^1 X^2 X^3 = 1$

Boosted black brane

$$ds^2 = -H^{2/3}(r) f(r) u_\mu u_\nu dx^\mu dx^\nu - 2 H^{-1/6}(r) u_\mu dx^\mu dr + r^2 H^{1/3}(r) (\eta_{\mu\nu} + u_\mu u_\nu) dx^\mu dx^\nu$$

$$A^a = \left(A_0^a(r) u_\mu + \mathcal{A}_\mu^a \right) dx^\mu$$

$$X^a = \frac{H^{1/3}(r)}{H_a(r)}$$

$$A_0^a(r) = \frac{\sqrt{m q^a}}{r^2 + q^a}$$

$$f(r) = -\frac{m}{r^2} + r^2 H(r)$$

$$H(r) = \prod_{a=1}^3 H^a(r)$$

$$H^a(r) = 1 + \frac{q^a}{r^2}$$

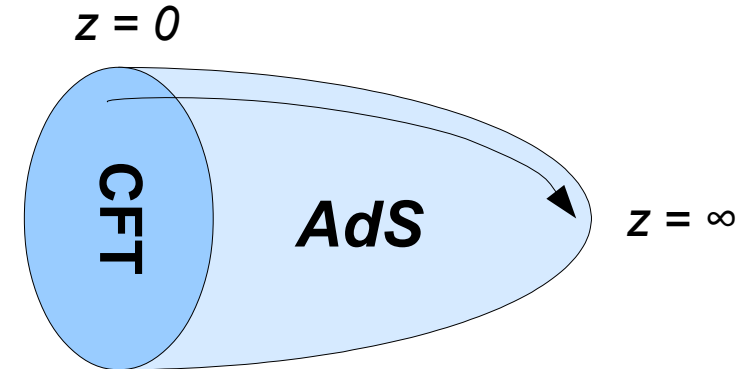
Torabian and Yee (2009)

Next order

We slowly vary 4-velocity and background fields

$$u_\mu = (-1, x^\nu \partial_\nu u_i)$$

$$\mathcal{A}_\mu^a = (0, x^\nu \partial_\nu \mathcal{A}_\mu^a)$$



Then solve equations of motion for this case and find corrections to the metric, gauge fields and scalars.

And consider the near-boundary expansion (Fefferman-Graham coordinates):

$$ds^2 = \frac{1}{z^2} \left(g_{\mu\nu}(z, x) dx^\mu dx^\nu + dz^2 \right),$$

$$T_{\mu\nu} = \frac{g_{\mu\nu}^{(4)}(x)}{4\pi G_5} + \dots$$

$$g_{\mu\nu}(z, x) = \eta_{\mu\nu} + g_{\mu\nu}^{(2)}(x) z^2 + g_{\mu\nu}^{(4)}(x) z^4 + \dots$$

$$A_\mu^a(z, x) = \mathcal{A}_\mu^a(x) + A_\mu^{a(2)}(x) z^2 + \dots$$

$$j_a^\mu = \frac{\eta^{\mu\nu} A_{a\nu}^{(2)}(x)}{8\pi G_5} + \dots$$

Transport coefficients

$$T^{\mu\nu} = \frac{m}{16\pi G_5} (\eta^{\mu\nu} + 4u^\mu u^\nu) + \dots,$$

$$j^{a\mu} = \frac{\sqrt{mq^a}}{8\pi G_5} u^\mu + \xi_\omega^a \frac{1}{2} \epsilon^{\mu\nu\lambda\rho} u_\nu \partial_\lambda u_\rho + \xi_B^{ab} \epsilon^{\mu\nu\lambda\rho} u_\nu \partial_\lambda A_\rho^b + \dots$$

(zeroth order)

$$\frac{\sqrt{mq^a}}{2m} = \frac{\rho^a}{\epsilon + P}$$

$$\xi_\omega^a = \frac{1}{16\pi G_5} \left(S^{abc} \mu^b \mu^c - \frac{\sqrt{mq^a}}{3m} S^{bcd} \mu^b \mu^c \mu^d \right)$$

$$\xi_B^{ab} = \frac{1}{16\pi G_5} \left(S^{abc} \mu^c - \frac{\sqrt{mq^a}}{4m} S^{bcd} \mu^c \mu^d \right)$$

(first order)

$$\mu^a \equiv A_0^a(r_H) - A_0^a(\infty)$$

$S_{abc} = 16\pi G_5 \cdot C_{abc}$ $\longrightarrow$ **We recover the hydrodynamic result!**

Thank you for the attention!

and

Have a good time!